

isibang/ms/2005/21

CONVOLUTION OPERATOR AND MAXIMAL FUNCTION FOR DUNKL TRANSFORM

SUNDARAM THANGAVELU AND YUAN XU

ABSTRACT. For a family of weight functions, h_{κ} , invariant under a finite reflection group on \mathbb{R}^d , analysis related to the Dunkl transform is carried out for the weighted L^p spaces. Making use of the generalized translation operator and the weighted convolution, we study the summability of the inverse Dunkl transform, including as examples the Poisson integrals and the Bochner-Riesz means. We also define a maximal function and use it to prove the almost everywhere convergence.

1. Introduction

The classical Fourier transform, initially defined on $L^1(\mathbb{R}^d)$, extends to an isometry of $L^2(\mathbb{R}^d)$ and it commutes with the rotation group. For a family of weight functions h_{κ} invariant under a reflection group G, there is a similar isometry of $L^2(\mathbb{R}^d, h_{\kappa}^2)$, called Dunkl transform ([3]), which enjoys properties similar to those of the classical Fourier transform. We denote this transform by \hat{f} in the following. It is defined by

$$\widehat{f}(x) = \int_{\mathbb{R}^d} E(x, -iy) f(y) h_{\kappa}^2(y) dy$$

where the usual character $e^{-i\langle x,y\rangle}$ is replaced by $E(x,-iy)=V_{\kappa}(e^{-i\langle \cdot,y\rangle})(x)$, in which V_{κ} is a positive linear operator (see the next section). If the parameter $\kappa=0$ then $h_{\kappa}(x)\equiv 1$ and $V_{\kappa}=\mathrm{id}$, so that \widehat{f} becomes the classical Fourier transform.

The basic properties of the Dunkl transform have been studied in [3, 7, 12, 14] and also in [11, 18] (see also the references therein). These studies are mostly for $L^2(\mathbb{R}^d)$ or for Schwartz class functions.

The purpose of this paper is to develop an L^p theory for the summability of the inverse Dunkl transform and prove a maximal inequality that implies almost everywhere convergence.

The classical Fourier transform behaves well with the translation operator $f \mapsto f(\cdot - y)$, which leaves the Lebesgue measure on \mathbb{R}^d invariant. However, the measure $h_{\kappa}^2(x)dx$ is no longer invariant under the usual translation. One ends up with a generalized translation operator, defined on the Dunkl transform side by

$$\widehat{\tau_y f}(x) = E(y, -ix)\widehat{f}(x), \qquad x \in \mathbb{R}^d.$$

An explicit formula for τ_y is unknown in general. In fact τ_y may not even be a positive operator. Consequently even the boundedness of τ_y in $L^p(\mathbb{R}^d; h^2_\kappa)$ becomes a challenging problem. At the moment an explicit formula for $\tau_y f$ is known only in two cases: when f is a radial function and when $G = \mathbb{Z}_2^d$. Properties of τ_y are studied in Section 3. In particular, the boundedness of the τ_y for radial functions is established.

Date: March 16, 2005.

 $^{1991\ \}textit{Mathematics Subject Classification}.\ 42\text{A}38,\ 42\text{B}08,\ 42\text{B}15.$

Key words and phrases. Dunkl transforms, reflection invariance, translation operator, convolution, summability.

ST wishes to thank YX for the warm hospitality during his stay in Eugene. The work of YX was supported in part by the National Science Foundation under Grant DMS-0201669.

For f, g in $L^2(\mathbb{R}^d; h^2_{\kappa})$ their convolution can be defined in terms of the translation operator as

$$(f *_{\kappa} g)(x) = \int_{\mathbb{R}^d} f(y) \tau_x g^{\vee}(y) h_{\kappa}^2(y) dy.$$

Based on a sharp Paley-Wiener theorem we are able to prove that $\phi_{\varepsilon} *_{\kappa} f$ converges to f in $L^{p}(\mathbb{R}^{d}; h_{\kappa}^{2})$ for certain radial ϕ , where ϕ_{ε} is a proper dilation of ϕ . This and other results are given in Section 4.

The convolution $*_{\kappa}$ can be used to study the summability of the inverse Dunkl transform. We prove the L^p convergence of the summability under mild conditions, including as examples Gaussian means (heat kernel transform), Abel means and the Bochner-Riesz means for the Dunkl transform in Section 5.

In the Section 6 we define a maximal function and prove that it is strong type (p,p) for 1 and weak type <math>(1,1). As usual, the maximal inequality implies almost everywhere convergence for the summability.

In the case $G = \mathbb{Z}_2^d$, the generalized translation operator is bounded in $L^p(\mathbb{R}^d; h_{\kappa}^2)$. Many of the results proved in the previous sections hold under conditions that are more relaxed in this case and the proof is more conventional. This case will be discussed in Section 7.

The following section is devoted to the preliminaries and background. The basic properties of the Dunkl transform will also be given.

2. Preliminaries

Let G be a finite reflection group on \mathbb{R}^d with a fixed positive root system R_+ , normalized so that $\langle v,v\rangle=2$ for all $v\in R_+$, where $\langle x,y\rangle$ denotes the usual Euclidean inner product. For a nonzero vector $v\in\mathbb{R}^d$, let σ_v denote the reflection with respect to the hyperplane perpendicular to $v, x\sigma_v:=x-2(\langle x,v\rangle/\|v\|^2)v, x\in\mathbb{R}^d$. Then G is a subgroup of the orthogonal group generated by the reflections $\{\sigma_v:v\in R_+\}$.

In [1], Dunkl defined a family of first order differential-difference operators, \mathcal{D}_i , that play the role of the usual partial differentiation for the reflection group structure. Let κ be a nonnegative multiplicity function $v \mapsto \kappa_v$ defined on R_+ with the property that $\kappa_u = \kappa_v$ whenever σ_u is conjugate to σ_v in G; then $v \mapsto \kappa_v$ is a G-invariant function. Dunkl's operators are defined by

$$\mathcal{D}_i f(x) = \partial_i f(x) + \sum_{v \in R_+} k_v \frac{f(x) - f(x\sigma_v)}{\langle x, v \rangle} \langle v, \varepsilon_i \rangle, \qquad 1 \le i \le d,$$

where $\varepsilon_1, \ldots, \varepsilon_d$ are the standard unit vectors of \mathbb{R}^d . These operators map \mathcal{P}_n^d to \mathcal{P}_{n-1}^d , where \mathcal{P}_n^d is the space of homogeneous polynomials of degree n in d variables. More importantly, these operators mutually commute; that is, $\mathcal{D}_i \mathcal{D}_j = \mathcal{D}_j \mathcal{D}_i$.

Associated with the reflection group and the function κ is the weight function h_{κ} defined by

(2.1)
$$h_{\kappa}(x) = \prod_{v \in R_{+}} |\langle x, v \rangle|^{\kappa_{v}}, \qquad x \in \mathbb{R}^{d}.$$

This is a positive homogeneous function of degree $\gamma_{\kappa} := \sum_{v \in R_{+}} \kappa_{v}$, and it is invariant under the reflection group G. The simplest example is given by the case $G = \mathbb{Z}_{2}^{d}$ for which h_{κ} is just the product weight function

$$h_{\kappa}(x) = \prod_{i=1}^{d} |x_i|^{\kappa_i}, \qquad \kappa_i \ge 0.$$

The Dunkl transform is taken with respect to the measure $h_{\kappa}^{2}(x)dx$.

There is a linear isomorphism that intertwines the algebra generated by Dunkl's operators with the algebra of partial differential operators. The intertwining operator V_{κ} is a linear operator determined uniquely by

$$V_{\kappa} \mathcal{P}_n \subset \mathcal{P}_n, \quad V_{\kappa} 1 = 1, \quad \mathcal{D}_i V_{\kappa} = V_{\kappa} \partial_i, \quad 1 \leq i \leq d.$$

The explicit formula of V_{κ} is not known in general. For the group $G = \mathbb{Z}_2^d$, it is an integral transform

$$(2.2) V_{\kappa}f(x) = b_{\kappa} \int_{[-1,1]^d} f(x_1t_1,\dots,x_dt_d) \prod_{i=1}^d (1+t_i)(1-t_i^2)^{\kappa_i-1} dt.$$

If some $\kappa_i = 0$, then the formula holds under the limit relation

$$\lim_{\lambda \to 0} b_{\lambda} \int_{-1}^{1} f(t)(1-t)^{\lambda-1} dt = [f(1) + f(-1)]/2.$$

It is known that V_{κ} is a positive operator ([12]); that is, $p \geq 0$ implies $V_{\kappa}p \geq 0$.

The function $E(x, y) := V_{\kappa}^{(x)} \left[e^{\langle x, y \rangle} \right]$, where the superscript means that V_{κ} is applied to the x variable, plays an important role in the development of the Dunkl transform. Some of its properties are listed below ([2]).

Proposition 2.1. For $x, y \in \mathbb{R}^n$,

- (1) E(x,y) = E(y,x);
- (2) $|E(x,y)| \le e^{\|x\| \cdot \|y\|}, \quad x, y \in \mathbb{C}^n;$
- (3) Let $\nu(z) = z_1^2 + \ldots + z_d^2$, $z_i \in \mathbb{C}$. For $z, w \in \mathbb{C}^d$,

$$c_h \int_{\mathbb{R}^d} E(z, x) E(w, x) h_{\kappa}^2(x) e^{-\|x\|^2/2} dx = e^{(\nu(z) + \nu(w))/2} E(z, w),$$

where c_h is the constant defined by $c_h^{-1} = \int_{\mathbb{R}^d} h_\kappa^2(x) e^{-\|x\|^2/2} dx$.

In particular, the function

$$E(x, iy) = V_{\kappa}^{(x)} \left[e^{i\langle x, y \rangle} \right], \qquad x, y \in \mathbb{R}^d,$$

plays the role of $e^{i\langle x,y\rangle}$ in the ordinary Fourier analysis. The Dunkl transform is defined in terms of it by

(2.3)
$$\widehat{f}(y) = c_h \int_{\mathbb{R}^d} f(x) E(x, -iy) h_{\kappa}^2(x) dx.$$

If $\kappa = 0$ then $V_{\kappa} = id$ and the Dunkl transform coincides with the usual Fourier transform. If d = 1 and $G = \mathbb{Z}_2$, then the Dunkl transform is related closely to the Hankel transform on the real line. In fact, in this case,

$$E(x, -iy) = \Gamma(\kappa + 1/2)(|xy|/2)^{-\kappa + 1/2} \left[J_{\kappa - 1/2}(|xy|) - i \operatorname{sign}(xy) J_{\kappa + 1/2}(|xy|) \right],$$

where J_{α} denotes the usual Bessel function

(2.4)
$$J_{\alpha}(t) = \left(\frac{t}{2}\right)^{\alpha} \sum_{n=0}^{\infty} \frac{(-1)^n}{n!\Gamma(n+\alpha+1)} \left(\frac{t}{2}\right)^{2n}.$$

We list some of the known properties of the Dunkl transform below ([3, 7]).

Proposition 2.2. (1) For $f \in L^1(\mathbb{R}^d; h_{\kappa}^2)$, \widehat{f} is in $C_0(\mathbb{R}^d)$.

(2) When both f and \hat{f} are in $L^1(\mathbb{R}^d; h^2_{\kappa})$ we have the inversion formula

$$f(x) = \int_{\mathbb{R}^d} E(ix, y) \widehat{f}(y) h_{\kappa}^2(y) dy.$$

- (3) The Dunkl transform extends to an isometry of $L^2(\mathbb{R}^d; h^2_{\kappa})$.
- (4) For Schwartz class functions f, $\widehat{\mathcal{D}_{i}f}(y) = iy_{i}\widehat{f}(y)$.

There are two more results that we will need. They require a little more preparation. First we need the definition of h-harmonics. The h-Laplacian is defined by $\Delta_h = \mathcal{D}_1^2 + \ldots + \mathcal{D}_d^2$ and it plays the role similar to that of the ordinary Laplacian. Let \mathcal{P}_n^d denote the subspace of homogeneous polynomials of degree n in d variables. An h-harmonic polynomial P of degree n is a homogeneous polynomial $P \in \mathcal{P}_n^d$ such that $\Delta_h P = 0$. Furthermore, let $\mathcal{H}_n^d(h_\kappa^2)$ denote the space of h-harmonic polynomials of degree n in d variables and define

$$\langle f, g \rangle_{\kappa} := a_{\kappa} \int_{S^{d-1}} f(x)g(x)h_{\kappa}^{2}(x)d\omega(x),$$

where $a_{\kappa}^{-1} = \int_{S^{d-1}} h_{\kappa}^2(x) d\omega$. Then $\langle P, Q \rangle_{\kappa} = 0$ for $P \in \mathcal{H}_n^d(h_{\kappa}^2)$ and $Q \in \Pi_{n-1}^d$. The spherical h-harmonics are the restriction of h-harmonics to the unit sphere. The standard Hilbert space theory shows that

$$L^{2}(h_{\kappa}^{2}) = \sum_{n=0}^{\infty} \bigoplus \mathcal{H}_{n}^{d}(h_{\kappa}^{2}).$$

Throughout this paper, we fix the value of $\lambda := \lambda_{\kappa}$ as

(2.5)
$$\lambda := \gamma_{\kappa} + \frac{d-2}{2} \quad \text{with} \quad \gamma_{\kappa} = \sum_{v \in R_{+}} \kappa_{v}.$$

Using the spherical-polar coordinates x = rx', where $x' \in S^{d-1}$, we have

(2.6)
$$\int_{\mathbb{R}^d} f(x) h_{\kappa}^2(x) dx = \int_0^{\infty} \int_{S^{d-1}} f(rx') h_{\kappa}^2(x') d\omega(x') r^{2\lambda_{\kappa} + 1} dr$$

from which it follows that

$$c_h^{-1} = \int_{\mathbb{R}^d} h_\kappa^2(x) e^{-\|x\|^2/2} dx = 2^{\lambda_\kappa} \Gamma(\lambda_\kappa + 1) a_\kappa^{-1}.$$

The following formula is useful for computing the Dunkl transform of certain functions ([3]).

Proposition 2.3. Let $f \in \mathcal{H}_n^d(h_{\kappa}^2)$, $y \in \mathbb{R}^d$ and $\mu > 0$. Then the function

$$g(x) = a_{\kappa} \int_{S^{d-1}} f(\xi) E(x, -i\mu\xi) h_{\kappa}^{2}(\xi) d\omega(\xi)$$

satisfies $\Delta_h g = -\mu^2 g$ and

$$g(x) = (-i)^n f\left(\frac{x}{\|x\|}\right) \left(\frac{\mu\|x\|}{2}\right)^{-\lambda_{\kappa}} J_{n+\lambda_{\kappa}}(\mu\|x\|).$$

We will also use the Hankel transform H_{α} defined on the positive reals \mathbb{R}_{+} . For $\alpha > -1/2$,

(2.7)
$$H_{\alpha}f(s) := \frac{1}{\Gamma(\alpha+1)} \int_0^{\infty} f(r) \frac{J_{\alpha}(rs)}{(rs)^{\alpha}} r^{2\alpha+1} dr.$$

The inverse Hankel transform is given by

(2.8)
$$f(r) = \frac{1}{\Gamma(\alpha+1)} \int_0^\infty H_\alpha f(s) \frac{J_\alpha(rs)}{(rs)^\alpha} s^{2\alpha+1} ds,$$

which holds under mild conditions on f; for example, it holds if f is piecewise continuous and of bounded variation in every finite subinterval of $(0, \infty)$, and $\sqrt{r}f \in L^1(\mathbb{R}_+)$ ([19, p. 456]).

Proposition 2.4. If $f(x) = f_0(||x||)$, then $\widehat{f}(x) = H_{\lambda_x} f_0(||x||)$.

Proof. This follows immediately from (2.6) and Proposition 2.3.

3. Generalized translation

One of the important tools in the classical Fourier analysis is the convolution

$$(f * g)(x) = \int_{\mathbb{R}^d} f(y)g(x - y)dy,$$

which depends on the translation $\tau_y: f(x) \mapsto f(x-y)$. There is a generalized translation for the reflection invariant weight function, which we study in this section.

3.1. Basic properties and explicit formulas. Taking the Fourier transform, we see that the translation $\tau_y f = f(\cdot - y)$ of \mathbb{R}^d satisfies $\widehat{\tau_y f}(x) = e^{-i\langle x,y\rangle} \widehat{f}(x)$. Looking at the Fourier transform side, an analogue of the translation operator for the Dunkl transform can be defined as follows:

Definition 3.1. Let $y \in \mathbb{R}^d$ be given. The generalized translation operator $f \mapsto \tau_y f$ is defined on $L^2(\mathbb{R}^d; h^2_{\kappa})$ by the equation

(3.1)
$$\widehat{\tau_y f}(x) = E(y, -ix)\widehat{f}(x), \qquad x \in \mathbb{R}^d.$$

Note that the definition makes sense as the Dunkl transform is an isometry of $L^2(\mathbb{R}^d; h_{\kappa}^2)$ onto itself and the function E(y, -ix) is bounded. When the function f is in the Schwartz class the above equation holds pointwise. Otherwise it is to be interpreted as an equation for L^2 functions. As an operator on $L^2(\mathbb{R}^d; h_{\kappa}^2)$, τ_y is bounded. A priori it is not at all clear whether the translation operator can be defined for L^p functions for p different from 2. One of the important issues is to prove the L^p boundedness of the translation operator on the dense subspace of Schwartz class functions. If it can be done then we can extend the definition to all L^p functions.

The above definition gives $\tau_y f$ as an L^2 function. It is useful to have a class of functions on which (3.1) holds pointwise. One such class is given by the subspace

$$A_{\kappa}(\mathbb{R}^d) = \{ f \in L^1(\mathbb{R}^d; h_{\kappa}^2) : \widehat{f} \in L^1(\mathbb{R}^d; h_{\kappa}^2) \}.$$

Note that $A_{\kappa}(\mathbb{R}^d)$ is contained in the intersection of $L^1(\mathbb{R}^d; h^2_{\kappa})$ and L^{∞} and hence is a subspace of $L^2(\mathbb{R}^d; h^2_{\kappa})$. For $f \in A_{\kappa}(\mathbb{R}^d)$ we have

(3.2)
$$\tau_y f(x) = \int_{\mathbb{R}^d} E(ix, \xi) E(-iy, \xi) \widehat{f}(\xi) h_{\kappa}^2(\xi) d\xi.$$

Before stating some properties of the generalized translation operator let us mention that there is an abstract formula for τ_y given in terms of the intertwining operator V_{κ} and its inverse. It takes the form of ([18])

(3.3)
$$\tau_y f(x) = V_{\kappa}^{(x)} \otimes V_{\kappa}^{(y)} \left[(V_{\kappa}^{-1} f)(x - y) \right]$$

for f being Schwartz class functions. We note that V_{κ}^{-1} satisfies the formula $V_{\kappa}^{-1}f(x) = e^{-\langle y, \mathcal{D} \rangle}f(x)|_{y=0}$. The above formula, however, does not provide much information on $\tau_y f$. The generalized translation operator has been studied in [12, 14, 18]. In [18] the equation (3.3) is taken as the starting point.

The following proposition collects some of the elementary properties of this operator which are easy to prove when both f and g are from $A_{\kappa}(\mathbb{R}^d)$.

Proposition 3.2. Assume that $f \in A_{\kappa}(\mathbb{R}^d)$ and $g \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ is bounded. Then

$$(1) \int_{\mathbb{R}^d} \tau_y f(\xi) g(\xi) h_{\kappa}^2(\xi) d\xi = \int_{\mathbb{R}^d} f(\xi) \tau_{-y} g(\xi) h_{\kappa}^2(\xi) d\xi.$$

(2)
$$\tau_{u} f(x) = \tau_{-x} f(-y)$$
.

Proof. The property (2) follows from the definition since $E(\lambda x, \xi) = E(x, \lambda \xi)$ for any $\lambda \in \mathbb{C}$. To prove (1) assume first that both f and g are from $A_{\kappa}(\mathbb{R}^d)$. Then both integrals in (1) are well defined. From the definition

$$\int_{\mathbb{R}^d} \tau_y f(\xi) g(\xi) h_{\kappa}^2(\xi) d\xi = \int_{\mathbb{R}^d} \left(\int_{\mathbb{R}^d} E(ix, \xi) E(-iy, \xi) \widehat{f}(\xi) h_{\kappa}^2(\xi) d\xi \right) g(x) h_{\kappa}^2(x) dx$$
$$= \int_{\mathbb{R}^d} \widehat{f}(\xi) \widehat{g}(-\xi) E(-iy, \xi) h_{\kappa}^2(\xi) d\xi.$$

We also have

$$\begin{split} \int_{\mathbb{R}^d} f(\xi) \tau_{-y} g(\xi) h_{\kappa}^2(\xi) d\xi &= \int_{\mathbb{R}^d} \left(\int_{\mathbb{R}^d} E(ix,\xi) E(iy,\xi) \widehat{g}(\xi) h_{\kappa}^2(\xi) d\xi \right) f(x) h_{\kappa}^2(x) dx \\ &= \int_{\mathbb{R}^d} \widehat{f}(-\xi) \widehat{g}(\xi) E(iy,\xi) h_{\kappa}^2(\xi) d\xi \\ &= \int_{\mathbb{R}^d} \widehat{f}(\xi) \widehat{g}(-\xi) E(-iy,\xi) h_{\kappa}^2(\xi) d\xi. \end{split}$$

This proves (1) when both f and g are from $A_{\kappa}(\mathbb{R}^d)$.

Suppose now $f \in A_{\kappa}(\mathbb{R}^d)$ but g is in the intersection of $L^1(\mathbb{R}^d; h_{\kappa}^2)$ and L^{∞} . Note that $g \in L^2(\mathbb{R}^d; h_{\kappa}^2)$ and so $\tau_y g$ is defined as an L^2 function. Since f is in $L^2(\mathbb{R}^d; h_{\kappa}^2)$ and bounded, both integrals are finite. The equation

$$\int_{\mathbb{R}^d} f(\xi)\widehat{g}(\xi)h_{\kappa}^2(\xi)d\xi = \int_{\mathbb{R}^d} \widehat{f}(\xi)g(\xi)h_{\kappa}^2(\xi)d\xi$$

which is true for Schwartz class functions remains true for $f,g\in L^2(\mathbb{R}^d;h^2_\kappa)$ as well. Using this we get

$$\int_{\mathbb{R}^d} \tau_y f(x) g(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} \tau_y f(-x) g(-x) h_{\kappa}^2(x) dx$$
$$= \int_{\mathbb{R}^d} E(y, -i\xi) \widehat{f}(\xi) \widehat{g}(-\xi) h_{\kappa}^2(\xi) d\xi.$$

By the same argument the integral on the right hand side is also given by the same expression. Hence \Box is proved.

We need to prove further properties of τ_y . In the classical case the ordinary translation satisfies

$$\int_{\mathbb{R}^d} f(x-y)dx = \int_{\mathbb{R}^d} f(x)dx.$$

Such a property is true for τ_y if f is a Schwartz class function. Indeed

$$\int_{\mathbb{R}^d} \tau_y f(x) h_{\kappa}^2(x) dx = \widehat{(\tau_y f)}(0) = \widehat{f}(0).$$

Here we have used the fact that τ_y takes S into itself. For $f \in A_{\kappa}(\mathbb{R}^d)$ though $\tau_y f$ is defined we do not know if it is integrable. We now address the question whether the above property holds at least for a subclass of functions.

For this purpose we make use of the following result which gives an explicit formula for $\tau_y f$ when f is radial, see [14].

Proposition 3.3. Let $f \in A_{\kappa}(\mathbb{R}^d)$ be radial and let $f(x) = f_0(||x||)$. Then

$$\tau_y f(x) = V_{\kappa} \left[f_0 \left(\sqrt{\|x\|^2 + \|y\|^2 - 2\|x\| \|y\| \langle x', \cdot \rangle} \right) \right] (y').$$

A special case of the above theorem is the following formula

(3.4)
$$\tau_y q_t(x) = e^{-t(\|x\|^2 + \|y\|^2)} E(2tx, y)$$

where

$$q_t(x) = (2t)^{-(\gamma + \frac{d}{2})} e^{-t||x||^2}$$

is the so called heat kernel. This formula has already appeared in [11]. The other known formula for $\tau_y f$ is the case when $G = \mathbb{Z}_2^d$.

Theorem 3.4. Let $f \in A_{\kappa}(\mathbb{R}^d)$ be radial and nonnegative. Then $\tau_y f \geq 0, \tau_y f \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ and

$$\int_{\mathbb{R}^d} \tau_y f(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} f(x) h_{\kappa}^2(x) dx.$$

Proof. As f is radial, the explicit formula in Proposition ?? shows that $\tau_y f \geq 0$ since V_{κ} is a positive operator. Taking $g(x) = e^{-t||x||^2}$ and making use of (3.4) we get

$$\int_{\mathbb{R}^d} \tau_y f(x) e^{-t\|x\|^2} h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} f(x) e^{-t(\|x\|^2 + \|y\|^2)} E(\sqrt{2t}x, \sqrt{2t}y) h_{\kappa}^2(x) dx.$$

As $|E(x,y)| \le e^{\|x\| \|y\|}$ we can take limit as $t \to 0$ to get

$$\lim_{t\to 0} \int_{\mathbb{R}^d} \tau_y f(x) e^{-t\|x\|^2} h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} f(x) h_{\kappa}^2(x) dx.$$

Since $\tau_y f \geq 0$, monotone convergence theorem applied to the integral on the left completes the proof. \Box

We would like to relax the condition on f in the above proposition. In order to do that we introduce the notion of generalized (Dunkl) convolution.

Definition 3.5. for $f, g \in L^2(\mathbb{R}^d; h^2_{\kappa})$ we define

$$f *_{\kappa} g(x) = \int_{\mathbb{R}^d} f(y) \tau_x g^{\vee}(y) h_{\kappa}^2(y) dy$$

where $g^{\vee}(y) = g(-y)$.

Note that as $\tau_x g^{\vee} \in L^2(\mathbb{R}^d; h_{\kappa}^2)$ the above convolution is well defined. We can also write the definition as

$$f *_{\kappa} g(x) = \int_{\mathbb{R}^d} \widehat{f}(\xi) \widehat{g}(\xi) E(ix, \xi) h_{\kappa}^2(\xi) d\xi.$$

If we assume that g is also in $L^1(\mathbb{R}^d; h^2_{\kappa})$ so that \widehat{g} is bounded, then by Plancherel theorem we obtain

$$||f *_{\kappa} g||_{\kappa,2} \le ||g||_{\kappa,1} ||f||_{\kappa,2}$$

We are interested in knowing under what conditions on g the operator $f \to f *_{\kappa} g$ defined on the Schwartz class can be extended to $L^p(\mathbb{R}^d; h_{\kappa}^2)$ as a bounded operator. But now we use the L^2 boundedness of the convolution to prove the following.

Theorem 3.6. Let $g \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ be radial, bounded and nonnegative. Then $\tau_y g \geq 0, \tau_y g \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ and

$$\int_{\mathbb{R}^d} \tau_y g(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} g(x) h_{\kappa}^2(x) dx.$$

Proof. Let q_t be the heat kernel defined earlier so that $\widehat{q}_t(\xi) = e^{-t\|\xi\|^2}$. By Plancherel theorem

$$||g *_{\kappa} q_t - g||_{\kappa,2}^2 = \int_{\mathbb{R}^d} |\widehat{g}(\xi)|^2 (1 - e^{-t||\xi||^2})^2 h_{\kappa}^2(\xi) d\xi$$

which shows that $g *_{\kappa} q_t \to g$ in $L^2(\mathbb{R}^d; h_{\kappa}^2)$ as $t \to 0$. Since τ_y is bounded on $L^2(\mathbb{R}^d; h_{\kappa}^2)$ we have $\tau_y(g *_{\kappa} q_t) \to \tau_y g$ in $L^2(\mathbb{R}^d; h_{\kappa}^2)$ as $t \to 0$. By passing to a subsequence if necessary we can assume that the convergence is also almost everywhere.

Now as g is radial and nonnegative, the convolution

$$g *_{\kappa} q_t(x) = \int_{\mathbb{R}^d} g(y) \tau_x q_t(y) h_{\kappa}^2(y) dy$$

is also radial and nonnegative. We also note that $g *_{\kappa} q_t \in A_{\kappa}(\mathbb{R}^d)$ as g is both in $L^1(\mathbb{R}^d; h_{\kappa}^2)$ and $L^2(\mathbb{R}^d; h_{\kappa}^2)$; in fact $g *_{\kappa} q_t \in L^1(\mathbb{R}^d, h_{\kappa}^2)$ as $q_t \in A_{\kappa}(\mathbb{R}^d)$ and, by Plancherel theorem and Hölder's inequality, $\|\widehat{g} *_{\kappa} q_t\|_{\kappa,1} = \|\widehat{g} \cdot \widehat{q}_t\|_{\kappa,1} \leq \|g\|_{\kappa,2} \|q_t\|_{\kappa,2}$. Thus by Theorem 3.6 we know that $\tau_y(g *_{\kappa} q_t)(x) \geq 0$. This gives us

$$\lim_{t \to 0} \tau_y(g *_{\kappa} q_t)(x) = \tau_y g(x) \ge 0$$

for almost every x. Once the nonnegativity of $\tau_y g(x)$ is proved it is easy to show that it is integrable. As before

$$\int_{\mathbb{R}^d} \tau_y g(x) e^{-t||x||^2} h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} g(x) e^{-t(||x||^2 + ||y||^2)} E(\sqrt{2t}x, \sqrt{2t}y) h_{\kappa}^2(x) dx.$$

Taking limit as t goes to 0 and using monotone convergence theorem we get

$$\int_{\mathbb{R}^d} \tau_y g(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} g(x) h_{\kappa}^2(x) dx.$$

This completes the proof.

We are now in a position to prove the following result. Let $L^p_{\text{rad}}(\mathbb{R}^d; h^2_{\kappa})$ denote the space of all radial functions in $L^p(\mathbb{R}^d; h^2_{\kappa})$.

Theorem 3.7. The generalized translation operator τ_y , initially defined on the intersection of $L^1(\mathbb{R}^d; h_{\kappa}^2)$ and L^{∞} , can be extended to all radial functions in $L^p(\mathbb{R}^d; h_{\kappa}^2)$, $1 \leq p \leq 2$, and $\tau_y : L^p_{\rm rad}(\mathbb{R}^d, h_{\kappa}^2) \to L^p(\mathbb{R}^d; h_{\kappa}^2)$ is a bounded operator.

Proof. Let $R: L^1(\mathbb{R}^d; h^2_{\kappa}) \to L^1(\mathbb{R}^d; h^2_{\kappa})$ be the radialisation defined by

$$Rf(x) = \int_{SO(n)} f(\sigma x) d\sigma.$$

Then for real valued $f \in L^1(\mathbb{R}^d; h_{\kappa}^2) \cap L^{\infty}$, Rf is radial and belongs to the intersection of $L^1(\mathbb{R}^d; h_{\kappa}^2)$ and L^{∞} . Since $-|Rf| \leq Rf \leq |Rf|$ the nonnegativity of τ_y on radial functions in $L^1(\mathbb{R}^d; h_{\kappa}^2) \cap L^{\infty}$ shows that $|\tau_y Rf(x)| \leq \tau_y |Rf|(x)$. Hence

$$\int_{\mathbb{R}^d} |\tau_y Rf(x)| h_{\kappa}^2(x) dx \leq \int_{\mathbb{R}^d} |Rf|(x) h_{\kappa}^2(x) dx \leq \|f\|_{\kappa,1}.$$

We also have $\|\tau_y Rf\|_{\kappa,2} \leq \|f\|_{\kappa,2}$. By interpolation we get $\|\tau_y Rf\|_{\kappa,p} \leq \|f\|_{\kappa,p}$ for all $1 \leq p \leq 2$. Thus the composition $\tau_y \circ R$ can be extended as a bounded operator on $L^p(\mathbb{R}^d; h^2_{\kappa})$. This proves the theorem.

Theorem 3.8. For every $f \in L^1_{\mathrm{rad}}(\mathbb{R}^d; h^2_{\kappa})$,

$$\int_{\mathbb{R}^d} \tau_y f(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} f(x) h_{\kappa}^2(x) dx.$$

Proof. Choose radial functions $f_n \in A_{\kappa}(\mathbb{R}^d)$ so that $f_n \to f$ and $\tau_y f_n \to \tau_y f$ in $L^1(\mathbb{R}^d; h^2_{\kappa})$. Since

$$\int_{\mathbb{R}^d} \tau_y f_n(x) g(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} f_n(x) \tau_{-y} g(x) h_{\kappa}^2(x) dx$$

for every $g \in A_{\kappa}(\mathbb{R}^d)$ we get, taking limit as n tends to infinity,

$$\int_{\mathbb{R}^d} \tau_y f(x) g(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} f(x) \tau_{-y} g(x) h_{\kappa}^2(x) dx.$$

Now take $g(x) = e^{-t||x||^2}$ and take limit as t goes to 0.. Since $\tau_y f \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ by dominated convergence theorem we obtain

$$\int_{\mathbb{R}^d} \tau_y f(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} f(x) h_{\kappa}^2(x) dx$$

for $f \in L^1(\mathbb{R}^d; h_{\kappa}^2)$.

We remark that it is still an open problem whether $\tau_y f$ can be defined for all $f \in L^1(\mathbb{R}^d; h^2_{\kappa})$.

3.2. **Positivity of** τ_y . As an immediate consequence of the explicit formula for the generalized translation of radial functions, if $f(x) \in A_{\kappa}(\mathbb{R}^d)$ is nonnegative, then $\tau_y f(x) \geq 0$ for all $y \in \mathbb{R}^d$ ([14]).

One would naturally expect that the generalized translation defines a positive operator; that is, $\tau_y f(x) \geq 0$ whenever $f(x) \geq 0$. This, however, turns out not to be the case. For $G = \mathbb{Z}_2$, the explicit formula given in Section 7 shows that τ_y is not positive in general (signed hypergroup, see [10]). Below we give an example to show that τ_y is not positive in a case where the explicit formula is not available. It depends on a method of computing generalized translation of simple functions. The explicit formula (3.2) can be used to define $\tau_y f$ when f is a polynomial.

Lemma 3.9. Let $y \in \mathbb{R}^d$. For $1 \le j \le d$, $\tau_y\{x_j\} = x_j - y_j$; and for $1 \le j, k \le d$,

$$\tau_y\{x_jx_k\} = (x_j - y_j)(x_k - y_k) - \kappa \sum_{v \in R_+} \left[V_\kappa(\langle x, y \rangle) - V_\kappa(\langle x\sigma_v, y \rangle) \right].$$

Proof. We use (3.4) and the fact that $\mathcal{D}_j \tau_y = \tau_y \mathcal{D}_j$. On the one hand, since the difference part of \mathcal{D}_j becomes zero when applied to radial functions,

$$\tau_y \mathcal{D}_j e^{-t||x||^2} = -2t\tau_y \left(\{\cdot\}_j e^{-t||\cdot||^2} \right) (x).$$

On the other hand it is easy to verify that

$$\mathcal{D}_j \tau_y e^{-t\|x\|^2} = \mathcal{D}_j \left[e^{-t(\|x\|^2 + \|y\|^2)} E(2tx,y) \right] = 2t e^{-t(\|x\|^2 + \|y\|^2)} E(2tx,y) (y_j - x_j).$$

Together, this leads to the equation

(3.5)
$$\tau_y \left(x_j e^{-t\|x\|^2} \right) = 2e^{-t(\|x\|^2 + \|y\|^2)} E(2tx, y) (x_j - y_j).$$

Taking the limit as $t \to 0$ gives $\tau_y\{x_j\} = x_j - y_j$.

Next we repeat the above argument, taking (3.5) as the starting point. Using the product formula for \mathcal{D}_k [5, p. 156], a simple computation gives

$$\begin{split} \mathcal{D}_k \tau_y \left(x_j e^{-t \|x\|^2} \right) &= \mathcal{D}_k \left[(x_j - y_j) e^{-t (\|x\|^2 + \|y\|^2)} \right] \\ &= e^{-t (\|x\|^2 + \|y\|^2)} \left[-2t (x_j - y_j) (x_k - y_k) E(2tx, y) \right. \\ &+ \delta_{k,j} E(2tx, y) + 2 \sum_{v \in R_+} \kappa_v \frac{v_k v_j}{\|v\|^2} E(2tx\sigma_v, y) \right]. \end{split}$$

On the other hand, computing $\mathcal{D}_k(x_i e^{-t||x||^2})$ leads to

$$\tau_y \left(\mathcal{D}_k(x_j e^{-t\|x\|^2}) \right) = -2t\tau_y(x_j x_k e^{-t\|x\|^2}) + \tau_y e^{-t\|x\|^2} \left[\delta_{k,j} + 2 \sum_{v \in R_+} \kappa_v \frac{v_k v_j}{\|v\|^2} \right].$$

Hence, using (3.4), the equation $\mathcal{D}_k \tau_y(x_j e^{-t||x||^2}) = \tau_y \mathcal{D}_k(x_j e^{-t||x||^2})$ gives

$$\begin{split} \tau_y \left(x_j x_k e^{-t \|x\|^2} \right) &= e^{-t (\|x\|^2 + \|y\|^2)} \Big[(x_j - y_j) (x_k - y_k) E(2tx, y) \\ &+ \sum_{v \in R_+} \kappa_v \frac{v_k v_j}{\|v\|^2} \frac{E(2tx, y) - E(2tx\sigma_v, y)}{t} \Big]. \end{split}$$

Taking the limit as $t \to 0$ gives the formula of $\tau_y\{x_jx_k\}$.

Proposition 3.10. The generalized translation τ_y is not a positive operator for the symmetric group S_d .

Proof. The formula $\tau_y\{x_jx_k\}$ depends on the values of $V_{\kappa}x_j$. For symmetric group S_d of d objects, the formula of $V_{\kappa}x_j$ is given by ([4])

$$V_{\kappa}x_j = \frac{1}{d\kappa + 1}(1 + x_j + \kappa|x|), \qquad |x| = x_1 + \ldots + x_d.$$

Let x(j,k) denote the transposition of x_j and x_k variables. It follows that

$$\begin{split} \tau_y \{x_j^2\} &= (x_j - y_j)^2 + \kappa \sum_{k \neq j} \left[V_{\kappa} (\langle x, y \rangle) - V_{\kappa} (\langle x(k, j), y \rangle) \right] \\ &= (x_j - y_j)^2 + \kappa \sum_{k \neq j} \left[(x_j - x_k) V_{\kappa} (y_j - y_k) \right] \\ &= (x_j - y_j)^2 + \frac{\kappa}{d\kappa + 1} \sum_{k \neq j} \left[(x_j - x_k) (y_j - y_k) \right]. \end{split}$$

Choosing $x = (1, 0, 0, \dots, 0)$ and $y = (0, 2, 2, \dots, 2)$, we see that $\tau_y(\{\cdot\}_1^2)(x) = -((d-2)\kappa+1)/(d\kappa+1) \le 0$. This proves the proposition.

Let us point out that, by (3.2), this proposition also shows that V_{κ}^{-1} is not a positive operator for the symmetric group. In the case of \mathbb{Z}_2 , an explicit formula of V_{κ}^{-1} is known ([21]) which is not positive.

3.3. Paley-Wiener theorem and the support of τ_y . In this subsection we prove a sharp Paley-Wiener theorem and study its consequences. Another type of Paley-Wiener theorem has been proved in [18]. A geometric form of the Paley-Wiener theorem is conjectured and studied recently in [8].

Let us denote by $\{Y_{j,n}: 1 \leq j \leq \dim \mathcal{H}_n^d(h_\kappa^2)\}$ an orthonormal basis of $\mathcal{H}_n^d(h_\kappa^2)$. First we prove a Paley-Wiener theorem for the Dunkl transform.

Theorem 3.11. Let $f \in$

 $SandBbeapositive number. Then fissupported in \{x: ||x|| \leq B\}$ if and only if for every j and n, the function

$$F_{j,n}(\rho) = \rho^{-n} \int_{S^{d-1}} \widehat{f}(\rho x) Y_{j,n}(x) h_{\kappa}^{2}(x) d\omega(x)$$

extends to an entire function of $\rho \in \mathbb{C}$ satisfying the estimate

$$|F_{i,n}(\rho)| \le c_{j,n} e^{B|\Im \rho||}.$$

Proof. By the definition of \hat{f} and Proposition 2.3,

$$\begin{split} \int_{S^{d-1}} \widehat{f}(\rho x) Y_{j,n}(x) h_{\kappa}^2(x) d\omega(x) \\ &= c \int_{\mathbb{R}^d} \int_{S^{d-1}} E(y, -i\rho x) Y_{j,n}(x) h_{\kappa}^2(x) d\omega(x) f(y) h_{\kappa}^2(y) dy \\ &= c \int_{\mathbb{R}^d} f(y) Y_{j,n}(y') \frac{J_{\lambda_k + n}(\rho || y ||)}{(\rho || y ||)^{\lambda_k}} h_{\kappa}^2(y) dy \\ &= c \int_0^{\infty} f_{j,n}(r) \frac{J_{\lambda_k + n}(r\rho)}{(r\rho)^{\lambda_k}} r^{2\lambda_{\kappa} + n + 1} dr, \end{split}$$

where c is a constant and

$$f_{j,n}(r) = r^{-n} \int_{S^{d-1}} f(ry') Y_{j,n}(y') h_{\kappa}^{2}(y') d\omega(y').$$

Thus, $F_{j,n}$ is the Hankel transform of order $\lambda_{\kappa} + n$ of the function $f_{j,n}(r)$. The theorem then follows from the Paley-Wiener theorem for the Hankel transform (see, for example, [6]).

Corollary 3.12. A function $f \in$

Sissupported in $\{x: ||x|| \leq B\}$ if and only if \widehat{f} extends to an entire function of $\zeta \in \mathbb{C}^d$ which satisfies

$$|\widehat{f}(\zeta)| \le c \, e^{B \|\Im \zeta\|}.$$

Proof. The direct part follows from the fact that $E(x, -i\zeta)$ is entire and $|E(x, -i\zeta)| \le c e^{||x|| \cdot ||\Im\zeta||}$. For the converse we look at

$$\int_{\mathbb{S}^{d-1}} \widehat{f}(\rho \xi) Y_{j,n}(\xi) h_{\kappa}^{2}(\xi) d\xi, \qquad \rho \in \mathbb{C}.$$

This is certainly entire and, from the proof of the previous theorem, has a zero of order n at the origin. Hence,

$$\rho^{-n} \int_{S^{d-1}} \widehat{f}(\rho \xi) Y_{j,n}(\xi) h_{\kappa}^{2}(\xi) d\xi$$

is an entire function of exponential type B, from which the converse follows from the theorem.

Proposition 3.13. Let $f \in$

Sbesupported in $\{x: ||x|| \leq B\}$. Then $\tau_y f$ is supported in $\{x: ||x|| \leq B + ||y||\}$.

Proof. Let $g(x) = \tau_y f(x)$. Then $\widehat{g}(\xi) = E(y, -i\xi) \widehat{f}(\xi)$ extends to \mathbb{C}^d as an entire function of type B + ||y||.

This property of τ_y has appeared in [18]. We note that the explicit formula for τ_y shows that the support set of τ_y given in Proposition 3.13 is sharp.

An important corollary in this regard is the following result.

Theorem 3.14. If $f \in C_0^{\infty}(\mathbb{R}^d)$ is supported in $||x|| \leq B$ then $||\tau_y f - f||_p \leq c||y||(B + ||y||)^{\frac{N}{p}}$ for $1 \leq p \leq \infty$. Consequently, $\lim_{y \to 0} ||\tau_y f - f||_{\kappa, p} = 0$.

Proof. From the definition we have

$$\tau_y f(x) - f(x) = \int_{\mathbb{R}^d} \left(E(y, -i\xi) - 1 \right) E(x, i\xi) \widehat{f}(\xi) h_{\kappa}^2(\xi) d\xi.$$

Since $|E(x, i\xi)| \le c||y|| \, ||\xi||$ we have the estimate

$$\|\tau_y f - f\|_{\infty} \le c\|y\| \int_{\mathbb{R}^d} \|\xi\| |\widehat{f}(\xi)| h_{\kappa}^2(\xi) d\xi.$$

As $\tau_y f$ is supported in $||x|| \leq (B + ||y||)$ we obtain

$$\|\tau_y f - f\|_p \le c\|y\|(B + \|y\|)^{\frac{N}{p}}$$

which goes to zero as y goes to zero.

4. The generalized convolution

4.1. Convolution. Recall that in section 3 we have defined the convolution $f *_{\kappa} g$ when $f, g \in L^2(\mathbb{R}^d; h_{\kappa}^2)$ by

$$(fst_\kappa g)(x)=\int_{\mathbb{R}^d}f(y) au_xg^ee(y)h_\kappa^2(y)dy.$$

This convolution has been considered in [11, 18]. It satisfies the following basic properties:

- (1) $\widehat{f *_{\kappa} g} = \widehat{f} \cdot \widehat{g};$
- (2) $f *_{\kappa} g = g *_{\kappa} f$.

We have also noted that the operator $f \to f *_{\kappa} g$ is bounded on $L^2(\mathbb{R}^d; h^2_{\kappa})$ provided \widehat{g} is bounded. We are interested in knowing under what conditions on g the operator $f \to f *_{\kappa} g$ can be extended to L^p as a bounded operator. If only the generalized translation operator can be extended as a bounded operator on $L^p(\mathbb{R}^d; h^2_{\kappa})$, then the convolution will satisfy the usual Young's inequality. At present we can only say something about convolution with radial functions.

Theorem 4.1. Let g be a bounded radial function in $L^1(\mathbb{R}^d; h^2_{\kappa})$. Then

$$f *_{\kappa} g(x) = \int_{\mathbb{R}^d} f(y) \tau_x g^{\vee}(y) h_{\kappa}^2(y) dy$$

initially defined on the intersection of $L^1(\mathbb{R}^d; h_{\kappa}^2)$ and $L^2(\mathbb{R}^d; h_{\kappa}^2)$ extends to all $L^p(\mathbb{R}^d; h_{\kappa}^2)$, $1 \leq p \leq \infty$ as a bounded operator. In particular,

Proof. For $g \in L^1(\mathbb{R}^d; h^2_{\kappa})$ which is bounded and radial we have $|\tau_y g| \leq \tau_y |g|$ which shows that

$$\int_{\mathbb{D}^d} |\tau_y g(x)| h_{\kappa}^2(x) dx \le \int_{\mathbb{D}^d} |g(x)| h_{\kappa}^2(x) dx.$$

Therefore,

$$\int_{\mathbb{D}^d} |f *_{\kappa} g(x)|^2 h_{\kappa}^2(x) dx \le ||f||_{\kappa,1} ||g||_{\kappa,1}.$$

We also have $||f *_{\kappa} g||_{\infty} \le ||f||_{\infty} ||g||_{\kappa,1}$. By interpolation we obtain $||f *_{\kappa} g||_{\kappa,p} \le ||g||_{\kappa,1} ||f||_{\kappa,p}$.

For $\phi \in L^1(\mathbb{R}^d; h^2_{\kappa})$ and $\varepsilon > 0$, we define the dilation ϕ_{ε} by

(4.2)
$$\phi_{\varepsilon}(x) = \varepsilon^{-(2\gamma_{\kappa} + d)} \phi(x/\varepsilon).$$

A change of variables shows that

$$\int_{\mathbb{R}^d} \phi_{\varepsilon}(x) h_{\kappa}^2(x) dx = \int_{\mathbb{R}^d} \phi(x) h_{\kappa}^2(x) dx, \quad \text{for all } \varepsilon > 0.$$

Theorem 4.2. Let $\phi \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ be a bounded radial function and assume that $c_h \int_{\mathbb{R}^d} \phi(x) h_{\kappa}^2(x) dx = 1$. Then for $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$, $1 \leq p < \infty$,

$$\lim_{\varepsilon \to 0} \|f *_{\kappa} \phi_{\varepsilon} - f\|_{\kappa, p} = 0.$$

Proof. For a given $\eta > 0$ we choose $g \in C_0^{\infty}$ such that $||g - f||_{\kappa,p} < \eta/3$. The triangle inequality and (4.1) lead to

$$||f *_{\kappa} \phi_{\varepsilon} - f||_{\kappa,p} \le \frac{2}{3} \eta + ||g *_{\kappa} \phi_{\varepsilon} - g||_{\kappa,p}$$

where we have used $||g - f||_{\kappa,p} < \eta/3$. Since ϕ is radial we can choose a radial function $\psi \in C_0^{\infty}$ such that

$$\|\phi - \psi\|_{\kappa,1} \le (12\|g\|_{\kappa,p})^{-1}\eta.$$

If we let $a = c_h \int_{\mathbb{R}^d} \psi(y) h_{\kappa}^2(y) dy$ then by the triangle inequality, (4.1) and (4.2),

$$\begin{aligned} \|g *_{\kappa} \phi_{\varepsilon} - g\|_{\kappa,p} &\leq \|g\|_{\kappa,p} \|\phi - \psi\|_{\kappa,1} + \|g *_{\kappa} \psi_{\varepsilon} - ag\|_{\kappa,p} + |a - 1| \|g\|_{\kappa,p} \\ &\leq \eta/6 + \|g *_{\kappa} \psi_{\varepsilon} - ag\|_{\kappa,p} \end{aligned}$$

since $||g||_{\kappa,p}||\phi-\psi||_{\kappa,1} \leq \frac{\eta}{12}$ and

$$|a-1| = \left| c_h \int_{\mathbb{R}^d} \left(\phi_{\varepsilon}(x) - \psi_{\varepsilon}(x) \right) h_{\kappa}^2(x) dx \right| \le (12 ||g||_{\kappa,p})^{-1} \eta.$$

Thus

$$||f *_{\kappa} \phi_{\varepsilon} - f||_{\kappa,p} \le \frac{5}{6} \eta + ||g *_{\kappa} \psi_{\varepsilon} - ag||_{\kappa,p}.$$

Hence it suffices to show that $||g *_{\kappa} \psi_{\varepsilon} - ag||_{\kappa,p} \leq \eta/6$.

But now $g \in A_{\kappa}(\mathbb{R}^d)$ and so

$$g *_{\kappa} \phi_{\varepsilon}(x) = \int_{\mathbb{R}^d} g(y) \tau_x \phi_{\varepsilon}^{\vee}(y) h_{\kappa}^2(y) dy = \int_{\mathbb{R}^d} \tau_{-x} g(y) \phi_{\varepsilon}(-y) h_{\kappa}^2(y) dy.$$

We also know that $\tau_{-x}g(y) = \tau_{-y}g(x)$ as $g \in C_0^{\infty}$. Therefore,

$$g *_{\kappa} \phi_{\varepsilon}(x) = \int_{\mathbb{R}^d} \tau_y g(x) \phi_{\varepsilon}(y) h_{\kappa}^2(y) dy.$$

In view of this

$$g *_{\kappa} \psi_{\varepsilon}(x) - ag(x) = \int_{\mathbb{D}^d} (\tau_y g(x) - g(x)) \psi_{\varepsilon}(y) h_{\kappa}^2(y) dy$$

which gives by Minkowski's integral inequality

$$\|g *_{\kappa} \psi_{\varepsilon} - ag\|_{\kappa,p} \le \int_{\mathbb{R}^d} \|\tau_y g - g\|_{\kappa,p} \psi_{\varepsilon}(y) |h_{\kappa}^2(y) dy.$$

If g is supported in $||x|| \leq B$ then the estimate in Theorem 3.14 gives

$$||g *_{\kappa} \psi_{\varepsilon} - ag||_{\kappa,p} \le c \int_{\mathbb{R}^{d}} ||y|| (B + ||y||)^{\frac{N}{p}} |\psi_{\varepsilon}(y)| h_{\kappa}^{2}(y) dy$$
$$\le c\varepsilon \int_{\mathbb{R}^{d}} ||y|| (B + ||\varepsilon y||)^{\frac{N}{p}} |\psi(y)| h_{\kappa}^{2}(y) dy$$

which can be made smaller that $\frac{\eta}{6}$ by choosing ε small. This completes the proof of the theorem. \Box

The explicit formula in the case of $G = \mathbb{Z}_2^d$ allows us to prove an analogous result without the assumption that ϕ is radial, see Section 7.

4.2. **Spherical means.** As in [9], we define the spherical mean operator on $A_{\kappa}(\mathbb{R}^d)$ by

$$S_r f(x) := a_{\kappa} \int_{S^{d-1}} \tau_{ry} f(x) h_{\kappa}^2(y) d\omega(y).$$

The generalized convolution of f with a radial function can be expressed in terms of the spherical means $S_r f$. In fact, if $f \in A_{\kappa}(\mathbb{R}^d)$ and $g(x) = g_0(||x||)$ is an integrable radial function then, using the spherical-polar coordinates,

$$(f *_{\kappa} g)(x) = c_h \int_{\mathbb{R}^d} \tau_y f(x) g(y) h_{\kappa}^2(y) dy$$

$$= c_h \int_0^{\infty} r^{2\lambda_{\kappa} + 1} g_0(r) \int_{S^{d-1}} \tau_{ry'} f(x) h_{\kappa}^2(y') dy' dr$$

$$= \frac{c_h}{a_{\kappa}} \int_0^{\infty} S_r f(x) g_0(r) r^{2\lambda_{\kappa} + 1} dr.$$

Proposition 4.3. Let $f \in A_{\kappa}(\mathbb{R}^d)$. Then for r > 0 and $\xi \in \mathbb{R}^d$,

(4.3)
$$\widehat{S_r f}(\xi) = \widehat{f}(\xi) \frac{J_{\lambda_{\kappa}}(r \|\xi\|)}{(r \|\xi\|/2)^{\lambda_{\kappa}}}.$$

Proof. Recalling the definition of $S_r f$ and changing the order of the integrals

$$\widehat{S_r f}(\xi) = a_{\kappa} \int_{S^{d-1}} c_h \int_{\mathbb{R}^d} \tau_{ry'} f(x) E(x, -i\xi) h_{\kappa}^2(x) dx h_{\kappa}^2(y) d\omega(y)$$

$$= a_{\kappa} \int_{S^{d-1}} \widehat{\tau_{ry} f}(\xi) h_{\kappa}^2(y) d\omega(y)$$

$$= \widehat{f}(\xi) a_{\kappa} \int_{S^{d-1}} E(ry, \xi) h_{\kappa}^2(y) d\omega(y),$$

so that the result follows from Proposition 2.3.

If the generalized translation operator is bounded on L^p , then we can prove that the same is true of the spherical mean operator S_r . However, for a certain range of p, the operator S_r can be shown to be bounded on L^p . The proof is based on analytic interpolation theorem of Stein. In order to apply this interpolation theorem, it is necessary that $S_r f$ is defined for all integrable simple functions f. When $f = \chi_E$ where E is a measurable subset of \mathbb{R}^d with finite measure, then $\tau_y f$ is defined as an L^2 function. From the definition of $\tau_y f$ it is clear that $y \to \tau_y f$ is a continuous function of y taking values in $L^2(\mathbb{R}^d; h_\kappa^2)$. Hence, $S_r f$ is defined as an L^2 function for any integrable simple function f. This also follows from Proposition 4.3 since all such simple functions are in $L^2(\mathbb{R}^d, h_\kappa^2)$ and the multiplier $J_{\lambda_\kappa}(r \|\xi\|)/(r \|\xi\|/2)^{\lambda_\kappa}$ is bounded.

We only know that the inequality $\|\tau_y f\|_{\kappa,p} \leq c\|f\|_{\kappa,p}$ holds either if f is a radial function and $1 \leq p \leq 2$ or if $G = \mathbb{Z}_2^d$. In both cases, the spherical mean operator $S_r f$ is bounded on $L^p(\mathbb{R}^d; h_{\kappa}^2)$. In general we do not know if the inequality $\|\tau_y f\|_{\kappa,p} \leq c\|f\|_{\kappa,p}$ holds for a generic function $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$. Without this assumption, however, we can prove the boundedness of S_r on certain range of L^p .

Theorem 4.4. The spherical mean operator is bounded on $L^p(\mathbb{R}^d; h^2_{\kappa})$ for all p satisfying

$$\frac{d+2\gamma_{\kappa}+1}{d+2\gamma_{\kappa}}$$

Proof. By (4.3), $S_r f$ is a multiplier operator with the multiplier given by the function $J_{(N-2)/2}(r||\xi||)(r||\xi||/2)^{-(N-2)/2}$, where $N = d + 2\gamma_{\kappa}$. We also note that convolution with χ_{B_r} , B_r being the ball of radius r centered at 0, corresponds to the multiplier $J_{N/2}(r||\xi||)(r||\xi||/2)^{-N/2}$. This suggests us that we consider the analytic family of operators, T_r^{ζ} defined as follows.

Given p with $(N+1)/N , choose <math>\delta > 0$ such that $1/p = N/(N+1) - \delta/2$ and consider

$$\widehat{T_r^{\zeta}f}(\xi) = \widehat{f}(\xi) \frac{J_{\frac{N+1}{2}(\zeta+\delta) - \frac{1}{2}}(r\|\xi\|)}{(r\|\xi\|/2)^{\frac{N+1}{2}(\zeta+\delta) - \frac{1}{2}}}.$$

The operator T_r^{ζ} is defined for all integrable simple functions f since $t^{-\alpha}J_{\alpha}(t)$ is bounded even for complex values of α as long as $\Re(\alpha) > \frac{-1}{2}$ (which follows from the Poisson integral representation of Bessel functions). Let $\zeta = t + i\eta$. Using the well known formula

$$J_{\mu+\nu+1}(r) = \frac{t^{\nu+1}}{2^{\nu}\Gamma(\nu+1)} \int_0^1 J_{\mu}(rs)s^{\mu+1} (1-s^2)^{\nu} ds$$

we can write $T_r^{\zeta} f$ as

$$T_r^{\zeta} f(x) = c \int_0^1 T_r^{\frac{N+1}{2}t - \frac{1}{2}} f(x) s^{(N+1)t} (1 - s^2)^{\frac{N+1}{2}(\delta + i\eta) - 1} ds.$$

When $t=0,\,T_r^{i\eta}$ is $L^2(\mathbb{R}^d;h_\kappa^2)$ bounded since it has a bounded multiplier. When $t=1,\,T_r^{1+i\eta}$ is given by

$$T_r^{1+i\eta}f(x) = c \int_0^1 (f *_{\kappa} \chi_{B_{rs}})(x) s^{(N+1)} (1-s^2)^{\frac{N+1}{2}(\delta+i\eta)-1} ds.$$

Hence, $T_r^{1+i\eta}$ is bounded on $L^p(\mathbb{R}^d; h_{\kappa}^2)$ for $1 \leq p \leq \infty$. Analytic interpolation ([17, p. 205]) between these two results gives us that T_r^t is bounded on $L^p(\mathbb{R}^d; h_{\kappa}^2)$ when $p^{-1} = (1+t)/2$. Choosing $t = (N-1)/(N+1) - \delta$ we see that $T_r^t f = S_r f$ which is bounded on $L^p(\mathbb{R}^d; h_{\kappa}^2)$ with $p^{-1} = N/(N+1) - \delta/2$ according to our choice of δ . This proves the theorem.

5. Summability of the inverse Dunkl Transform

Let $\Phi \in L^1(\mathbb{R}^d; h^2_{\kappa})$ and assume $\Phi(0) = 1$. For $f \in \mathcal{S}$ and $\varepsilon > 0$ define

$$T_{\varepsilon}f(x) = c_h \int_{\mathbb{R}^d} \widehat{f}(y) E(ix, y) \Phi(-\varepsilon y) h_{\kappa}^2(y) dy$$

We study the convergence of $T_{\varepsilon}f$ as $\varepsilon \to 0$. Note that $T_0f = f$ by the inversion formula for the Dunkl transform. If $T_{\varepsilon}f$ can be extended to all $f \in L^p(\mathbb{R}^d; h^2_{\kappa})$ and if $T_{\varepsilon}f \to f$ in $L^p(\mathbb{R}^d; h^2_{\kappa})$, we say that the inverse Dunkl transform is Φ -summable.

Proposition 5.1. Let Φ and $\phi = \widehat{\Phi}$ both belong to $L^1(\mathbb{R}^d; h^2_{\kappa})$. If Φ is radial then

$$T_{\varepsilon} f(x) = (f *_{\kappa} \phi_{\varepsilon})(x)$$

for all $f \in L^2(\mathbb{R}^d; h_{\kappa}^2)$ and $\varepsilon > 0$.

Proof. Under the hypothesis on Φ both T_{ε} and the operator taking f into $(f *_{\kappa} \phi_{\varepsilon})$ extend to $L^{2}(\mathbb{R}^{d}; h_{\kappa}^{2})$ as bounded operators. So it is enough to verify $T_{\varepsilon}f(x) = (f *_{\kappa} \phi_{\varepsilon})(x)$ for all f in the Schwartz class. By the definition of the Dunkl transform,

$$T_{\varepsilon}f(x) = c_{h} \int_{\mathbb{R}^{d}} \widehat{\tau_{-x}f}(y) \Phi(-\varepsilon y) h_{\kappa}^{2}(y) dy$$

$$= c_{h} \int_{\mathbb{R}^{d}} \tau_{-x}f(\xi) c_{h} \int_{\mathbb{R}^{d}} \Phi(-\varepsilon y) E(y, -i\xi) h_{\kappa}^{2}(y) dy h_{\kappa}^{2}(\xi) d\xi$$

$$= c_{h} \varepsilon^{-(d+2\gamma_{\kappa})} \int_{\mathbb{R}^{d}} \tau_{-x}f(\xi) \widehat{\Phi}(-\varepsilon^{-1}\xi) h_{\kappa}^{2}(\xi) d\xi$$

$$= (f *_{\kappa} \phi_{\varepsilon})(x)$$

where we have changed variable $\xi \mapsto -\xi$ and used the fact that $\tau_{-x}f(-\xi) = \tau_{\xi}f(x)$.

If the radial function ϕ satisfies the conditions of Theorem 4.2 we obtain the following result.

Theorem 5.2. Let $\Phi(x) \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ be radial and assume that $\widehat{\Phi} \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ is bounded and $\Phi(0) = 1$. For $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$, $T_{\varepsilon}f$ converges to f in $L^p(\mathbb{R}^d; h_{\kappa}^2)$ as $\varepsilon \to 0$, for $1 \le p < \infty$.

We consider several examples. In our first example we take Φ to be the Gaussian function, $\Phi(x) = e^{-\|x\|^2/2}$. By (3) of Proposition (2.1) with z = iy and w = 0, $\widehat{\Phi}(x) = e^{-\|x\|^2/2}$. We choose $\varepsilon = 1/\sqrt{2t}$ and define

$$q_t(x) = \Phi_{\varepsilon}(x) = (2t)^{-(\gamma_k + \frac{d}{2})} e^{-\|x\|^2/4t}.$$

Then $q_t(x)$ satisfies the heat equation for the h-Laplacian,

$$\Delta_h u(x,t) = \partial_t u(x,t),$$

where Δ_h is applied to x variables. For this Φ , our summability method is just $f *_{\kappa} q_t$. By (3.4), the generalized translation of q_t is given explicitly by

$$\tau_y q_t(x) = (2t)^{-(\gamma_k + \frac{d}{2})} e^{-(\|x\|^2 + \|y\|^2)/4t} E\left(\frac{x}{\sqrt{2t}}, \frac{y}{\sqrt{2t}}\right)$$

which is the heat kernel for the solution of the heat equation for h-Laplacian. Then a corollary of Theorem 5.2 gives the following result in [13].

Theorem 5.3. Suppose $f \in L^p(\mathbb{R}^d; h^2_{\kappa}), 1 \leq p < \infty$.

(1) The heat transform

$$H_t f(x) := (f *_{\kappa} q_t)(x) = c_h \int_{\mathbb{R}^d} f(y) \tau_y q_t(x) h_{\kappa}^2(y) dy, \qquad t > 0,$$

converges to f in $L^p(\mathbb{R}^d; h_{\kappa}^2)$ as $t \to 0$.

(2) Define $H_0f(x) = f(x)$. Then the function $H_tf(x)$ solves the initial value problem

$$\Delta_h u(x,t) = \partial_t u(x,t), \qquad u(x,0) = f(x), \qquad (x,t) \in \mathbb{R}^d \times [0,\infty).$$

Our second example is the analogue of the Poisson summability, where we take $\Phi(x) = e^{-\|x\|}$. This case has been studied in [15]. In this case, one can compute the Dunkl transform $\widehat{\Phi}$ just as in the case of the ordinary Fourier transform, namely, using

(5.1)
$$e^{-t} = \frac{1}{\sqrt{\pi}} \int_0^\infty \frac{e^{-u}}{\sqrt{u}} e^{-t^2/4u} du,$$

and making use of the fact that the transform of Gaussian is itself (see [17, p. 6]). The result is

$$\widehat{e^{-\|x\|}} = c_{d,\kappa} \frac{1}{(1 + \|x\|^2)^{\gamma_{\kappa} + \frac{d+1}{2}}}, \qquad c_{d,\kappa} = 2^{\gamma_{\kappa} + \frac{d}{2}} \frac{\Gamma(\gamma_{\kappa} + \frac{d+1}{2})}{\sqrt{\pi}}.$$

In this case, we define the Poisson kernel as the dilation of $\widehat{\Phi}$,

(5.2)
$$P_{\varepsilon}(x) := c_{d,\kappa} \frac{\varepsilon}{(\varepsilon^2 + ||x||^2)^{\gamma_{\kappa} + \frac{d+1}{2}}}.$$

Since $\Phi(0) = 1$, it is easy to see that $\int P(x,\varepsilon)h_{\kappa}^{2}(x)dx = 1$. We have

Theorem 5.4. Suppose $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$, $1 \leq p < \infty$. Then the Poisson integral $f *_{\kappa} P_{\varepsilon}$ converges to f in $L^p(\mathbb{R}^d; h_{\kappa}^2)$.

Again the proof is a corollary of Theorem 5.2. For $\kappa = 0$, it becomes the Poisson summability for the classical Fourier analysis on \mathbb{R}^d .

Next we consider the analogue of the Bochner-Riesz means for which

$$\Phi(x) = \begin{cases} (1 - ||x||^2)^{\delta}, & ||x|| \le 1, \\ 0, & \text{otherwise} \end{cases}$$

where $\delta > 0$. As in the case of the ordinary Fourier transform, we take $\varepsilon = 1/R$ where R > 0. Then the Bochner-Riesz means is defined by

$$S_R^{\delta} f(x) = c_h \int_{\|y\| \le R} \left(1 - \frac{\|y\|^2}{R^2} \right)^{\delta} \widehat{f}(y) E(ix, y) h_{\kappa}^2(y) dy.$$

Recall that we have defined $\lambda_{\kappa} = \frac{(d-2)}{2} + \gamma_{\kappa}$ and $N = d + 2\gamma_{\kappa}$.

Theorem 5.5. If $f \in L^p(\mathbb{R}^d; h^2_{\kappa})$, $1 and <math>\delta > (N-1)/2$, then

$$||S_R^{\delta} f - f||_{\kappa, p} \to 0, \quad as \ R \to \infty.$$

Proof. The proof follows as in the case of ordinary Fourier transform [17, p. 171]. From Proposition 2.4 and the properties of the Bessel function, we have

$$\widehat{\Phi}(x) = 2^{\lambda_{\kappa}} ||x||^{-\lambda_{\kappa} - \delta - 1} J_{\lambda_{\kappa} + \delta + 1}(||x||).$$

Hence, by $J_{\alpha}(r) = \mathcal{O}(r^{-1/2})$, $\widehat{\Phi} \in L^{1}(\mathbb{R}^{d}; h_{\kappa}^{2})$ under the condition $\delta > \lambda_{\kappa} + 1/2 = (N-1)/2$.

We note that $\lambda_{\kappa} = (d-2)/2 + \gamma_{\kappa}$ where γ_{κ} is the sum of all (nonnegative) parameters in the weight function. If all parameters are zero, then $h_{\kappa}(x) \equiv 1$ and we are back to the classical Fourier transform, for which the index (d-1)/2 is the critical index for the Bochner-Riesz means. We do not know if the index (N-1)/2 is the critical index for the Bochner-Riesz means of the Dunkl transforms.

6. Maximal function and almost everywhere summability

For $f \in L^2(\mathbb{R}^d; h^2_{\kappa})$ we define the maximal function $M_{\kappa}f$ by

$$M_{\kappa}f(x) = \sup_{r>0} \frac{1}{d_{\kappa}r^{d+2\gamma_{\kappa}}} |f *_{\kappa} \chi_{B_r}(x)|,$$

where χ_{B_r} is the characteristic function of the ball B_r of radius r centered at 0 and $d_{\kappa} = a_{\kappa}/(d+2\gamma_{\kappa})$. Using (2.6) we have $\int_{B_r} h_{\kappa}^2(y) dy = (a_{\kappa}/(d+2\gamma_{\kappa})) r^{d+2\gamma_{\kappa}}$. Therefore, we can also write $M_{\kappa} f(x)$ as

$$M_{\kappa}f(x) = \sup_{r>0} \frac{\left| \int_{\mathbb{R}^d} f(y) \tau_x \chi_{B_r}(y) h_{\kappa}^2(y) dy \right|}{\int_{B_r} h_{\kappa}^2(y) dy}.$$

If $\varphi \in C_0^{\infty}(\mathbb{R}^d)$ is a radial function such that $\chi_{B_r}(x) \leq \varphi(x)$ then from Theorem 3.6 it follows that $\tau_y \chi_{B_r}(x) \leq \tau_y \varphi(x)$. But $\tau_y \varphi$ is bounded; hence $\tau_y \chi_{B_r}$ is bounded and compactly supported so that it belongs to $L^p(\mathbb{R}^d; h_{\kappa}^2)$. This means that the maximal function $M_{\kappa}f$ is defined for all $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$. We also note that as $\tau_y \chi_{B_r} \geq 0$ we have $M_{\kappa}f(x) \leq M_{\kappa}|f|(x)$.

Theorem 6.1. The maximal function is bounded on $L^p(\mathbb{R}^d; h_{\kappa}^2)$ for 1 ; moreover it is of weak type <math>(1,1), that is, for $f \in L^1(\mathbb{R}^d; h_{\kappa}^2)$ and a > 0,

$$\int_{E(a)} h_{\kappa}^{2}(x) dx \le \frac{c}{a} ||f||_{\kappa,1}$$

where $E(a) = \{x : M_{\kappa}f(x) > a\}$ and c is a constant independent of a and f.

Proof. Without loss of generality we can assume that $f \ge 0$. Let $\sigma = d + 2\gamma_{\kappa} + 1$ and define for $j \ge 0$, $B_{r,j} = \{x : 2^{-j-1}r \le ||x|| \le 2^{-j}r\}$. Then

$$\chi_{B_{r,j}}(y) = (2^{-j}r)^{\sigma} (2^{-j}r)^{-\sigma} \chi_{B_{r,j}}(y)$$

$$\leq \left(\frac{4}{5}\right)^{\sigma} (2^{-j}r)^{\sigma-1} \frac{2^{-j}r}{((2^{-j}r)^2 + ||y||^2)^{\sigma/2}} \chi_{B_{r,j}}(y)$$

$$\leq c(2^{-j}r)^{\sigma-1} P_{2^{-j}r}(y),$$

where P_{ε} is the Poisson kernel defined in (5.2) and $c = c_{d,\kappa}(4/5)^{\sigma}$. Since χ_{B_r} and P_{ε} are both bounded, integrable radial functions, it follows from Theorem 3.6 that

$$\tau_x \chi_{B_{r,j}}(y) \le c(2^{-j}r)^{\sigma-1} \tau_x P_{2^{-j}r}(y).$$

This shows that for any positive integer m

$$\int_{\mathbb{R}^d} f(y) \sum_{j=0}^m \tau_x \chi_{B_{r,j}}(y) h_{\kappa}^2(y) dy \le c \sum_{j=0}^{\infty} (2^{-j}r)^{\sigma-1} \int_{\mathbb{R}^d} f(y) \tau_x P_{2^{-j}r}(y) h_{\kappa}^2(y) dy$$

$$\le c r^{d+2\gamma_{\kappa}} \sup_{t>0} f *_{\kappa} P_t(x).$$

As $\sum_{j=0}^{m} \chi_{B_{r,j}}(y)$ converges to $\chi_{B_r}(y)$ in $L^1(\mathbb{R}^d; h_{\kappa}^2)$, the boundedness of τ_x on $L^1_{\mathrm{rad}}(\mathbb{R}^d; h_{\kappa}^2)$ shows that $\sum_{j=0}^{m} \tau_x \chi_{B_{r,j}}(y)$ converges to $\tau_x \chi_{B_r}(y)$ in $L^1(\mathbb{R}^d; h_{\kappa}^2)$. By passing to a subsequence if necessary we can assume that $\sum_{j=0}^{m} \tau_x \chi_{B_{r,j}}(y)$ converges to $\tau_x \chi_{B_r}(y)$ for almost every y. Thus all the functions involved are uniformly bounded by $\tau_x \chi_{B_r}(y)$. This shows that $\sum_{j=0}^{m} \tau_x \chi_{B_{r,j}}(y)$ converges to $\tau_x \chi_{B_r}(y)$ in $L^{p'}(\mathbb{R}^d; h_{\kappa}^2)$ and hence

$$\lim_{m\to\infty} \int_{\mathbb{R}^d} f(y) \sum_{j=0}^m \tau_x \chi_{B_{r,j}}(y) h_{\kappa}^2(y) dy = \int_{\mathbb{R}^d} f(y) \tau_x \chi_{B_r}(y) h_{\kappa}^2(y) dy.$$

Thus we have proved that

$$f *_{\kappa} \chi_{B_r}(x) \le cr^{d+2\gamma_{\kappa}} \sup_{t>0} f *_{\kappa} P_t(x)$$

which gives the inequality $M_{\kappa}f(x) \leq cP^*f(x)$, where $P^*f(x) = \sup_{t>0} f *_{\kappa} P_t(x)$ is the maximal function associated to the Poisson semigroup.

Therefore it is enough to prove the boundedness of P^*f . Here we follow a general procedure used in [16]. By looking at the Dunkl transforms of the Poisson kernel and the heat kernel we infer that

$$f *_{\kappa} P_t(x) = \frac{r}{\sqrt{2\pi}} \int_0^{\infty} (f *_{\kappa} q_s)(x) e^{-r^2/2s} s^{-3/2} ds,$$

which implies, as in [16, p. 49], that

$$P^*f(x) \le \sup_{t>0} \frac{1}{t} \int_0^t Q_s f(x) ds,$$

where $Q_s f(x) = f *_{\kappa} q_s(x)$ is the heat semigroup. Hence using the Hopf- Dunford-Schwartz ergodic theorem as in [16, p. 48], we get the boundedness of P^*f on $L^p(\mathbb{R}^d; h_{\kappa}^2)$ for 1 and the weak type <math>(1,1).

The maximal function can be used to study almost everywhere convergence of $f *_{\kappa} \varphi_{\epsilon}$ as they can be controlled by the Hardy-Littlewood maximal function $M_{\kappa}f$ under some conditions on φ . Recall that $N = d + 2\gamma_{\kappa}$.

Theorem 6.2. Let $\phi \in A_{\kappa}(\mathbb{R}^d)$ be a real valued radial function which satisfies $|\phi(x)| \leq c(1+||x||)^{-N-1}$. Then

$$\sup_{\varepsilon>0}|f*_{\kappa}\phi_{\varepsilon}(x)|\leq cM_{\kappa}f(x).$$

Consequently, $f *_{\kappa} \phi_{\varepsilon}(x) \to f(x)$ for almost every x as ε goes to 0 for all f in $L^p(\mathbb{R}^d; h^2_{\kappa})$, $1 \le p < \infty$.

Proof. We can assume that both f and ϕ are nonnegative. Writing

$$\phi_{arepsilon}(y) = \sum_{j=-\infty}^{\infty} \phi_{arepsilon}(y) \chi_{arepsilon 2^{j} \leq \|y\| \leq arepsilon 2^{j+1}}(y)$$

we have

$$\sum_{j=-m}^{m} \phi_{\varepsilon}(y) \tau_{x} \chi_{\varepsilon^{2^{j}} \leq ||y|| \leq \varepsilon^{2^{j+1}}}(y) \leq c \sum_{j=-m}^{m} (1 + \varepsilon^{2^{j}})^{-N-1} \tau_{x} \chi_{\varepsilon^{2^{j}} \leq ||y|| \leq \varepsilon^{2^{j+1}}}(y).$$

This shows that

$$\int_{\mathbb{R}^d} f(y)\phi_{\varepsilon}(y) \sum_{j=-m}^m \chi_{\varepsilon 2^j \le ||y|| \le \varepsilon 2^{j+1}}(y) h_{\kappa}^2(y) dy \le c \sum_{j=-m}^m (1+\varepsilon 2^j)^{-N-1} (\varepsilon 2^j)^N M_{\kappa} f(x)$$

$$\le c M_{\kappa} f(x).$$

Since $\phi(y) \leq c(1+||y||)^{-N-1} \leq cP_1(y)$ it follows that $\tau_x\phi(y) \leq c\tau_xP_1(y)$ is bounded. Arguing as in the previous theorem we can show that the left hand side of the above inequality converges to $f *_{\kappa} \phi_{\varepsilon}(x)$. Thus we obtain

$$\sup_{\varepsilon>0}|f*_{\kappa}\phi_{\varepsilon}(x)|\leq cM_{\kappa}f(x)$$

from which the proof of almost everywhere convergence follows from the standard argument. \Box

The above two theorems show that the maximal functions $M_{\kappa}f$ and P^*f are comparable. As a corollary we obtain almost everywhere convergence of Bochner-Riesz means.

Corollary 6.3. When $\delta \geq \frac{N+1}{2}$ the Bochner-Riesz means $S_R^{\delta}f(x)$ converges to f(x) for almost every x for all $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$, $1 \leq p < \infty$.

We expect the corollary to be true for all $\delta > \frac{(N-1)}{2}$ as in the case of the Fourier transform. This can be proved if in the above theorem the hypothesis on ϕ can be relaxed to $|\phi(x)| \leq c(1+||x||)^{-N-\epsilon}$ for some $\epsilon > 0$. Since we do not know if $\tau_y((1+||x||)^{-N-1})$ is bounded or not we cannot repeat the proof of the above theorem.

7. Product weight function invariant under \mathbb{Z}_2^d

Recall that in the case $G = \mathbb{Z}_2^d$, the weight function h_{κ} is a product function

(7.1)
$$h_{\kappa}(x) = \prod_{i=1}^{d} |x_i|^{\kappa_i}, \qquad \kappa_i \ge 0.$$

In this case the explicit formula of the intertwining operator V_{κ} is known (see (2.2)) and there is an explicit formula for τ_y . The following formula is contained in [10], where it is studied under the context of signed hypergroups.

Theorem 7.1. For $G = \mathbb{Z}_2^d$ and h_{κ} in (7.1),

$$\tau_y f(x) = \tau_{y_1} \cdots \tau_{y_d} f(x), \qquad y = (y_1, \dots, y_d) \in \mathbb{R}^d,$$

where for $G = \mathbb{Z}_2$ and $h_{\kappa}(t) = |t|^{\kappa}$ on \mathbb{R} ,

(7.2)
$$\tau_{s}f(t) = \frac{1}{2} \int_{-1}^{1} f\left(\sqrt{t^{2} + s^{2} - 2stu}\right) \left(1 + \frac{t - s}{\sqrt{t^{2} + s^{2} - 2stu}}\right) \Phi_{\kappa}(u) du + \frac{1}{2} \int_{-1}^{1} f\left(-\sqrt{t^{2} + s^{2} - 2stu}\right) \left(1 - \frac{t - s}{\sqrt{t^{2} + s^{2} - 2stu}}\right) \Phi_{\kappa}(u) du,$$

where $\Phi_{\kappa}(u) = b_{\kappa}(1+u)(1-u^2)^{\kappa-1}$. Consequently, for each $y \in \mathbb{R}^d$, the generalized translation operator τ_y for \mathbb{Z}_2^d extends to a bounded operator on $L^p(\mathbb{R}^d; h_{\kappa}^2)$. More precisely, $\|\tau_y f\|_{\kappa,p} \leq 3\|f\|_{\kappa,p}$, $1 \leq p \leq \infty$.

Since the generalized translation operator τ_y extends to a bounded operator on $L^p(\mathbb{R}^d; h_{\kappa}^2)$, many results stated in the previous sections can be improved and the proofs can be carried out more conveniently as in the classical Fourier analysis. In particular, the properties of τ_y given in Proposition 3.2, Theorem 3.6 and Theorem 3.8 all hold under the more relaxed condition of $f \in L^1(\mathbb{R}^d; h_{\kappa}^2)$.

The standard proof [22] can now be used to show that the generalized convolution satisfies the following analogous of Young's inequality.

Proposition 7.2. Let $G = \mathbb{Z}_2^d$. Let $p, q, r \geq 1$ and $p^{-1} = q^{-1} + r^{-1} - 1$. Assume $f \in L^q(\mathbb{R}^d, h_{\kappa}^2)$ and $g \in L^r(\mathbb{R}^d, h_{\kappa}^2)$, respectively. Then

$$||f *_{\kappa} g||_{\kappa,p} \le c||f||_{\kappa,q}||g||_{\kappa,r}$$

In the following we give several results that improve the corresponding results in the previous sections significantly. We start with an improved version of Theorem 4.2. The boundedness of τ_y allows us to remove the assumption that ϕ is radial.

Theorem 7.3. Let $\phi \in L^1(\mathbb{R}^d, h_{\kappa}^2)$ and assume $\int_{\mathbb{R}^d} \phi(x) h_{\kappa}^2(x) dx = 1$. Then for $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$, $1 \leq p < \infty$, or $f \in C_0(\mathbb{R}^d)$ if $p = \infty$,

$$\lim_{\varepsilon \to 0} \|f *_{\kappa} \phi_{\varepsilon} - f\|_{\kappa, p} = 0, \qquad 1 \le p \le \infty.$$

Proof. First we assume that $f \in C_0^{\infty}(\mathbb{R}^d)$. By Theorem 3.14, $\|\tau_y f(x) - f(x)\|_{\kappa,p} \to 0$ as $y \to 0$ for $1 \le p \le \infty$. In general, for $f \in L^p(\mathbb{R}^d; h_{\kappa}^2)$ we write $f = f_1 + f_2$ where f_1 is continuous with compact support and $\|f_2\|_{\kappa,p} \le \delta$. Then the first term of the inequality

$$\|\tau_{u}f(x)-f(x)\|_{\kappa,p} \leq \|\tau_{u}f_{1}(x)-f_{1}(x)\|_{\kappa,p} + \|\tau_{u}f_{2}(x)-f_{2}(x)\|_{\kappa,p}$$

goes to zero as $\varepsilon \mapsto 0$ and the second term is bounded by $(1+c)\delta$ as $\|\tau_y f_2\|_{\kappa,p} \le c\|f\|_{\kappa,p}$. This proves that $\|\tau_y f(x) - f(x)\|_{\kappa,p} \to 0$ as $y \to 0$. We have then

$$\begin{split} c_h \int_{\mathbb{R}^d} |f *_{\kappa} g_{\varepsilon}(x) - f(x)|^p h_{\kappa}^2(x) dx \\ &= c_h \int_{\mathbb{R}^d} \left| c_h \int_{\mathbb{R}^d} (\tau_y f(x) - f(x)) g_{\varepsilon}(y) h_{\kappa}^2(y) dy \right|^p h_{\kappa}^2(x) dx \\ &\leq c_h \int_{\mathbb{R}^d} \|\tau_y f - f\|_{\kappa, p}^p |g_{\varepsilon}(x)| h_{\kappa}^2(x) dx \\ &= c_h \int_{\mathbb{R}^d} \|\tau_{\varepsilon y} f - f\|_{\kappa, p}^p |g(x)| h_{\kappa}^2(x) dx, \end{split}$$

which goes to zero as $\varepsilon \to 0$.

Our next result is about the boundedness of the spherical means operator.

Theorem 7.4. Let $G = \mathbb{Z}_2^d$. For $f \in L^p(\mathbb{R}^d, h_{\kappa}^2)$,

$$||S_r f||_{\kappa,p} \le c||f||_{\kappa,p}, \qquad 1 \le p \le \infty.$$

Furthermore, $||S_r f - f||_{\kappa,p} \mapsto 0$ as $r \to 0+$.

Proof. Using Hölder's inequality,

$$|S_r f(x)|^p \le a_\kappa \int_{S^{d-1}} | au_{ry} f(x)|^p h_\kappa^2(y) d\omega(y).$$

Hence, a simple computation shows that

$$\begin{aligned} c_h \int_{\mathbb{R}^d} |S_r f(x)|^p h_{\kappa}^2(x) dx &\leq c_h \int_{\mathbb{R}^d} a_{\kappa} \int_{S^{d-1}} |\tau_{ry} f(x)|^p h_{\kappa}^2(y) d\omega(y) h_{\kappa}^2(x) dx \\ &= a_{\kappa} \int_{S^{d-1}} \|\tau_{ry} f\|_{\kappa,p}^p h_{\kappa}^2(y) d\omega(y) \\ &\leq c \|f\|_{\kappa,p}. \end{aligned}$$

Furthermore, we have

$$||S_r f - f||_{\kappa,p}^p \le a_\kappa \int_{S^{d-1}} ||\tau_{ry} f - f||_{\kappa,p}^p h_\kappa^2(y) d\omega(y)$$

which goes to zero as $r \to 0$ since $\|\tau_{ry}f - f\|_{\kappa,p} \to 0$.

The boundedness of $\tau_y f$ in $L^p(\mathbb{R}^d; h^2_{\kappa})$ also allows us to relax the condition of Theorem 6.2.

Theorem 7.5. Set $G = \mathbb{Z}_2^d$. Let $\phi(x) = \phi_0(||x||) \in L^1(\mathbb{R}^d, h_{\kappa}^2)$ be a radial function. Assume that ϕ_0 is differentiable, $\lim_{r\to\infty} \phi_0(r) = 0$ and $\int_0^{\infty} r^{2\lambda_{\kappa}+2} |\phi_0(r)| dr < \infty$, then

$$|(f *_{\kappa} \phi)(x)| \le cM_{\kappa}f(x).$$

In particular, if $\phi \in L^1(\mathbb{R}^d, h_{\kappa}^2)$ and $c_h \int_{\mathbb{R}^d} \phi(x) h_{\kappa}^2(x) dx = 1$, then

- (1) For $1 \leq p \leq \infty$, $f *_{\kappa} \phi_{\varepsilon}$ converges to f as $\varepsilon \to 0$ in $L^{p}(\mathbb{R}^{d}; h_{\kappa}^{2})$;
- (2) For $f \in L^1(\mathbb{R}^d, h_{\kappa}^2)$, $(f *_{\kappa} \phi_{\varepsilon})(x)$ converges to f(x) as $\varepsilon \to 0$ for almost all $x \in \mathbb{R}^d$.

Proof. By definition of the spherical means $S_t f$, we can also write

$$M_{\kappa}f(x) = \sup_{r>0} \frac{\left| \int_0^r t^{2\lambda_{\kappa}+1} S_t f(x) dt \right|}{\int_0^r t^{2\lambda_{\kappa}+1} dt}.$$

Since $|M_{\kappa}f(x)| \leq cM_{\kappa}|f|(x)$, we can assume $f(x) \geq 0$. The assumption on ϕ_0 shows that

$$\lim_{r \to \infty} \phi_0(r) \int_0^r S_t f(x) t^{2\lambda_{\kappa} + 1} dt = \lim_{r \to \infty} \phi_0(r) \int_{\mathbb{R}^d} \tau_y f(x) h_{\kappa}^2(y) dy$$
$$= \lim_{r \to \infty} \phi_0(r) \int_{\mathbb{R}^d} f(y) h_{\kappa}^2(y) dy = 0.$$

Hence, using the spherical-polar coordinates and integrating by parts, we get

$$(f *_{\kappa} \phi)(x) = \int_0^{\infty} \phi_0(r) r^{2\lambda_{\kappa}+1} S_r f(x) dr$$
$$= -\int_0^{\infty} \int_0^r S_t f(x) t^{2\lambda_{\kappa}+1} dt \phi'(r) dr,$$

which implies that

$$|(f*_{\kappa}\phi)(x)| \le cM_{\kappa}f(x)\int_0^{\infty} r^{2\lambda_{\kappa}+2}|\phi_0'(r)|dr,$$

the boundedness of the last integral proves the maximal inequality.

As an immediate consequence of the this theorem, the Bochner-Riesz means converge almost everywhere if $\delta > (N-1)/2$ for $G = \mathbb{Z}_2^d$, which closes the gap left open in Corollary 6.3.

We can further enhance Theorem 7.5 by removing the assumption that ϕ is radial. For this purpose, we make the following simple observation about the maximal function: If f is nonnegative then we can drop the absolute value sign in the definition of the maximal function, even though $\tau_y f$ may not be nonnegative.

Lemma 7.6. If $f \in L^1(\mathbb{R}^d, h^2_{\kappa})$ is a nonnegative function then

$$M_{\kappa}f(x) = \sup_{r>0} \frac{\int_{B_r} \tau_y f(x) h_{\kappa}^2(y) dy}{\int_{B_r} h_{\kappa}^2(y) dy}.$$

In particular, if f and g are two nonnegative functions then

$$M_{\kappa}f + M_{\kappa}g = M_{\kappa}(f+g).$$

Proof. Since $\tau_y \chi_{B_r}(x)$ is nonnegative, we have

$$(f *_{\kappa} \chi_{B_r})(x) = \int_{\mathbb{R}^d} f(y) \tau_y \chi_{B_r}(x) h_{\kappa}^2(y) dy$$

is nonnegative if f is nonnegative. Hence we can drop the absolute value sign in the definition of $M_{\kappa}f$.

Theorem 7.7. Set $G = \mathbb{Z}_2^d$. Let $\phi \in L^1(\mathbb{R}^d, h_{\kappa}^2)$ and let $\psi(x) = \psi_0(||x||) \in L^1(\mathbb{R}^d, h_{\kappa}^2)$ be a nonnegative radial function such that $|\phi(x)| \leq \psi(x)$. Assume that ψ_0 is differentiable, $\lim_{r \to \infty} \psi_0(r) = 0$ and $\int_0^\infty r^{2\lambda_{\kappa}+2} |\psi_0(r)| dr < \infty$. Then $\sup_{\epsilon>0} |f*_{\kappa}\phi_{\epsilon}(x)|$ is of weak type (1,1). In particular, if $\phi \in L^1(\mathbb{R}^d, h_{\kappa}^2)$ and $c_h \int_{\mathbb{R}^d} \phi(x) h_{\kappa}^2(x) dx = 1$, then for $f \in L^1(\mathbb{R}^d, h_{\kappa}^2)$, $(f*_{\kappa}\phi_{\epsilon})(x)$ converges to f(x) as $\epsilon \to 0$ for almost all $x \in \mathbb{R}^d$.

Proof. Since $M_{\kappa}f(x) \leq M_{\kappa}|f|(x)$, we can assume that $f(x) \geq 0$. The proof uses the explicit formula for $\tau_y f$. Let us first consider the case of d=1. Since ψ is an even function, $\tau_y \psi$ is given by the formula

$$au_y f(x) = \int_{-1}^1 f\left(\sqrt{x^2 + y^2 - 2xyt}\right) \Phi_{\kappa}(t) dt$$

by (7.2). Since (x - y)(1 + t) = (x - yt) - (y - xt), we have

$$\frac{|x-y|}{\sqrt{x^2 + y^2 - 2xyt}}(1+t) \le 2.$$

Consequently, by the explicit formula of $\tau_y f$ (7.2), the inequality $|\phi(x)| \leq \psi(x)$ implies that,

$$|\tau_u \phi(x)| \le \tau_u \psi(x) + 2\widetilde{\tau}_u \psi(x),$$

where $\tilde{\tau}_y \psi$ is defined by

$$\widetilde{\tau}_y \psi(x) = b_{\kappa} \int_{-1}^1 f\left(\sqrt{x^2 + y^2 - 2xyt}\right) (1 - t^2)^{\kappa - 1} dt.$$

Note that $\tilde{\tau}_y \psi$ differs from $\tau_y \psi$ by a factor of (1+t) in the weight function. Changing variables $t \mapsto -t$ and $y \mapsto -y$ in the integrals shows that

$$\int_{\mathbb{R}} f(y)\widetilde{\tau}_y \psi(x) h_{\kappa}^2(y) dy = \int_{\mathbb{R}} F(y) \tau_y \psi(x) h_{\kappa}^2(y) dy.$$

where F(y) = (f(y) + f(-y))/2. Hence, it follows that

$$|(f*_{\kappa}\phi)(x)| = \left| \int_{\mathbb{D}} f(y)\tau_y \phi(x) h_{\kappa}^2(y) dy \right| \le (f*_{\kappa}\psi)(x) + 2(F*_{\kappa}\psi)(x).$$

The same consideration can be extended to the case of \mathbb{Z}_2^d for d > 1. Let $\{e_1, \ldots, e_d\}$ be the standard Euclidean basis. For $\delta_j = \pm 1$ define $x\delta_j = x - (1 + \delta_j)x_je_j$ (that is, multiplying the j-th component of x by δ_j gives $x\delta_j$). For $1 \leq j \leq d$ we define

$$F_{j_1,\dots,j_k} = 2^{-k} \sum_{(\delta_{j_1},\dots,\delta_{j_k}) \in \mathbb{Z}_2^k} f(x\delta_{j_1} \cdots \delta_{j_k}).$$

In particular, $F_j(x) = (F(x) + F(x\delta_j))/2$, $F_{j_1,j_2}(x) = (F(x) + F(x\delta_{j_1}) + F(x\delta_{j_2}) + F(x\delta_{j_2}))/4$, and the last sum is over \mathbb{Z}_2^d , $F_{1,\dots,d}(x) = 2^{-d} \sum_{\sigma \in \mathbb{Z}_2^d} f(x\sigma)$. Following the proof in the case of d = 1 it is not hard to see that

$$|(f *_{\kappa} \phi)(x)| \le (f *_{\kappa} \psi)(x) + 2 \sum_{j=1}^{d} (F_{j} *_{\kappa} \psi)(x) + 4 \sum_{j_{1} \neq j_{2}} (F_{j_{1}, j_{2}} *_{\kappa} \psi)(x) + \dots + 2^{d} (F_{1, \dots, d} *_{\kappa} \psi)(x).$$

For $G = \mathbb{Z}_2^d$, the explicit formula of τ_y shows that $M_{\kappa}f(x)$ is even in each of its variables. Hence, applying the result of the previous theorem on each of the above terms, we get

$$|(f *_{\kappa} \phi)(x)| \le M_{\kappa} f(x) + 2 \sum_{j=1}^{d} M_{\kappa} F_{j}(x) + 4 \sum_{j_1 \ne j_2} M_{\kappa} F_{j_1, j_2}(x) + \dots + 2^{d} M_{\kappa} F_{1, \dots, d}(x).$$

Since all F_j are clearly nonnegative, by Lemma 7.6, the last expression can be written as $M_{\kappa}H$, where H is the sum of all functions involved. Consequently, since $||F_{j_1,\ldots,j_d}||_{\kappa,1} \leq ||f||_{\kappa,1}$, it follows that

$$\int_{\{x:(f*_\kappa\phi)(x)\geq a\}} h_\kappa^2(y) dy \leq c \frac{\|H\|_{\kappa,1}}{a} \leq c_d \frac{\|f\|_{\kappa,1}}{a}.$$

Hence, $f *_{\kappa} \phi$ is of weak type (1,1), from which the almost everywhere convergence follows as usual. \square

We note that we do not know if the inequality $|(f *_{\kappa} \phi)(x)| \leq c M_{\kappa} f(x)$ holds in this case, since we only know $M_{\kappa}(R(\delta)f)(x) = R(\delta)M_{\kappa}f(x) = M_{\kappa}f(x\delta)$, where $R(\delta)f(x) = f(x\delta)$ for $\delta \in G$, from which we cannot deduce that $M_{\kappa}F_{j_1,\ldots,j_k}(x) \leq c M_{\kappa}f(x)$.

REFERENCES

- [1] C. F. Dunkl, Differential-difference operators associated to reflection groups, Trans. Amer. Math. Soc. 311 (1989),
- [2] C. F. Dunkl, Integral kernels with reflection group invariance, Canad. J. Math. 43 (1991), 1213-1227.
- [3] C. F. Dunkl, Hankel transforms associated to finite reflection groups, in Hypergeometric functions on domains of positivity, Jack polynomials, and applications (Tampa, FL, 1991), 123-138, Contemporary Mathematics 138, American Math. Society, Providence, RI. 1992.
- [4] C. F. Dunkl, Intertwining operators and polynomials associated with the symmetric group, *Monatsh. Math.* 126 (1998), 181-209.
- [5] C. F. Dunkl and Yuan Xu, Orthogonal polynomials of several variables, Cambridge Univ. Press, 2001.
- [6] T. H. Koornwinder, A new proof of a Paley-Wiener theorem for the Jacobi transform, Ark. Mat. 13 (1975), 145-159.
- [7] M. F. E. de Jeu, The Dunkl transform, Invent. Math. 113 (1993), 147-162.
- [8] M. F. E. de Jeu, Paley-Wiener theorems for the Dunkl transform, Preprint, April, 2004. ArXiv:math.CA/0404439.
- [9] H. Mejjaoli and K. Trimèche, On a mean value property associated with the Dunkl Laplacian operator and applications, *Integral Transform. Spec. Funct.* **12** (2001), 279-302.
- [10] M. Rösler, Bessel-type signed hypergroups on ℝ, in: H. Heyer, A. Mukherjea (eds.), Probability measures on groups and related structures XI, Proc. Oberwolfach 1994, World Scientific, Singapore, 1995, 292-304.
- [11] M. Rösler, Generalized Hermite polynomials and the heat equation for Dunkl operators, Comm. Math. Phys. 192, (1998) 519-542.
- [12] M. Rösler, Positivity of Dunkl's intertwining operator, Duke Math. J., 98 (1999), 445-463.
- [13] M. Rösler, One-parameter semigroups related to abstract quantum models of Calogero type, in: Infinite Dimensional Harmonic Analysis (Kyoto, Sept. 1999, eds. H. Heyer et al.) Grbner-Verlag, 2000, p. 290 - 305.
- [14] M. Rösler, A positive radial product formula for the Dunkl kernel, Trans. Amer. Math. Soc., 355 (2003), 2413-2438.
- [15] M. Rösler and M. Voit, Markov processes associated with Dunkl operators, Adv. Appl. Math. 21 (1998), 575-643.
- [16] E. M. Stein, *Topics in Harmonic Analysis related to the Littlewood-Paley theory*, Annals of Mathematical Studies 63, Princeton Univ. Press, Princeton, NJ, 1970.

- [17] E. M. Stein and G. Weiss, Introduction to Fourier analysis on Euclidean spaces, Princeton Univ. Press, Princeton, NJ, 1971.
- [18] K. Trimèche, Paley-Wiener theorems for the Dunkl transform and Dunkl translation operators, *Integral Transforms* and Special Functions, 13 (2002), 17-38.
- [19] G. N. Watson, A Treatise on the Theory of Bessel Functions, 2nd edition, Cambridge University Press, London, 1962.
- [20] Yuan Xu, Integration of the intertwining operator for h-harmonic polynomials associated to reflection groups, Proc. Amer. Math. Soc. 125 (1997), 2963-2973.
- [21] Yuan Xu, An integral formula for generalized Gegenbauer polynomials and Jacobi polynomials, *Adv. in Applied Math.* **29** (2002), 328-343.
- [22] A. Zygmund, Trigonometric Series, Cambridge University Press, 1959.

STAT-MATH DIVISION, INDIAN STATISTICAL INSTITUTE, 8TH MILE, MYSORE ROAD, BANGALORE-560 059, INDIA.

 $E ext{-}mail\ address: weluma@isibang.ac.in}$

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF OREGON, EUGENE, OREGON 97403-1222.

 $E\text{-}mail\ address: \verb"yuan@math.uoregon.edu"$