# Percolation theory: from classical to dependent models

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# Outline

Origin and basic definitions

Some questions and few results

Going beyond independence

More results and open questions

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#### What is Percolation?

# What is the likelihood that a liquid, say water, will pass through a porous medium?



(Image source: https://www.theseptictankstore.co.uk/blog/soil-percolation-porosity-test/)

#### What is Percolation?

#### Percolation model arose as a simple stochastic model for such a

#### situation (Broadbent and Hammersley'57)



(Image source: https://www.theseptictankstore.co.uk/blog/soil-percolation-porosity-test/)



Consider a connected graph  $\mathcal{G}$ , e.g. the square lattice  $\mathbb{Z}^2$  or a tree



Let p be a parameter in [0, 1] (the so-called density)



Declare every site to be open with probability p and closed otherwise



This model is called site percolation



In a different version (bond percolation) we open or close edges

#### How does it look like?



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#### We encode a percolation configuration by

$$\omega = (\omega_{oldsymbol{
u}}: oldsymbol{
u} \in \mathcal{G}) \in \{0,1\}^{\mathcal{G}}$$

where

$$\omega_{\mathbf{v}} = \begin{cases} 1 & \text{if } \mathbf{v} \text{ is open} \\ \\ 0 & \text{if } \mathbf{v} \text{ is closed} \end{cases}$$

The probability measure on the space  $\Omega$  of percolation configurations for density *p* is given by

$${\sf P}_p = \prod_{v\in \mathcal{G}} ~(p\delta_1 + (1-p)\delta_0)$$



Primary events of interest:

 $\{S \leftrightarrow T\} = \{S, T \subset G \text{ are connected by an open path}\}$ 

From now onwards  $\mathcal{G}$  will be an infinite connected graph

Notice that  $\mathcal{G}$  is a metric space w.r.t. the minimum path length

 $\Lambda_{n,x}$  is the ball of radius *n* around *x* in graph distance  $d_{\mathcal{G}}$ , i.e.

$$\Lambda_{n,x} = \{y \in \mathcal{G} : d_{\mathcal{G}}(x,y) \leq n\}$$

The boundary of  $\Lambda_n$  is  $\partial \Lambda_{n,x} = \Lambda_{n,x} \setminus \Lambda_{n-1,x}$ 

#### The first question: does it percolate?

The quantity to look at is the one-arm probability:

$$\theta_{n,x}(p) = \mathbf{P}_p\left[\overbrace{x}^{\partial \wedge_{n,x}}\right] = \mathbf{P}_p[x \leftrightarrow \partial \wedge_{n,x}]$$

Notice that  $\theta_{n,x}(p) \searrow \theta_x(p) = \mathbf{P}_p[x \leftrightarrow \infty]$  as  $n \to \infty$ 

Also notice that  $\theta_x(p) > 0$  if and only if  $\theta_y(p) > 0$  for all  $y \in \mathcal{G}$ 

#### Can we compute $\theta_{n,x}(p)$ explicitly? Let's make an attempt!

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#### Can we compute $\theta_{n,x}(p)$ explicitly? Let's make an attempt!

$$heta_{n,x}(
ho) = \sum_{\omega \in \{ 0 \longleftrightarrow \partial \Lambda_{n,x} \}} \mathbf{P}_{
ho}[\omega]$$

$$\mathbf{P}_{\rho}[\omega] = p^{\sum_{y \in \Lambda_{n,x}} \omega_{y}} (1-\rho)^{\sum_{y \in \Lambda_{n,x}} 1-\omega_{y}} \quad (\mathsf{Easy!})$$

Evaluate the sum over admissible configurations (VERY difficult!)

One can compare this with the difficulty of computing the partition function for models in statistical physics (e.g. the lsing model)

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In fact this analogy is far from being artificial!

Percolation provides one of the simplest yet extremely rich example of phase transition (e.g. solid-liquid-gas, ferromagnet-paramagnet)

# The function $\theta_x(p)$ : the phase diagram

The following properties are not difficult to see from the definition:

 $\theta_x(p)$  is non-decreasing in  $p, \ \theta(0) = 0$  and  $\theta(1) = 1$ 

## The function $\theta_x(p)$ : the phase diagram

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 $\theta_x(p)$  is non-decreasing in  $p, \ \theta(0) = 0$  and  $\theta(1) = 1$ 



#### The function $\theta_{x}(p)$ : the critical density

Therefore there exists a critical parameter  $p_c = p_c(\mathcal{G})$  defined as:

$$p_c = \sup\{p \in [0,1] : \theta(p) = 0\}$$



#### The function $\theta_{x}(p)$ : existence of phase transition

Notice that  $p_c$  can be a priori 0 or 1 (no phase transition)

$$p_c = \sup\{p \in [0,1] : \theta(p) = 0\}$$



#### Existence of phase transition

 $p_c$  is always positive! In fact,  $p_c(\mathcal{G}) \geq \frac{1}{\max_{v \in \mathcal{G}} \deg(v)}$  (Peierls'36)

The proof is essentially an energy-entropy type argument

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$$p_c(\mathbb{Z}^d) < 1$$
 for all  $d \geq 2$  whereas  $p_c(\mathbb{Z}) = 1$ 

Compare with the fact that the Ising model has no phase transition in dimension 1! (Ising'25)

Deriving a generic condition on  $\mathcal{G}$  ensuring  $p_c(\mathcal{G}) < 1$  is non-trivial

A natural guess would be that some "suitable" notion of dimension of G is strictly bigger than 1 (Benjamini-Schramm'96)

We say G has isoperimetric dimension at least d > 0 if

$$|\partial K| \geq c|K|^{rac{d-1}{d}}$$
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Theorem (Duminil-Copin, G., Raoufi, Severo and Yadin 2018)

Let G a graph of bounded degree with isoperimetric dimension strictly larger than 4, then  $p_c(G) < 1$ .

G is called quasi-transitive if the action of the automorphism group Aut(G) on G has finitely many orbits

Typical examples include Cayley graphs of finitely generated group

We say that  $\mathcal{G}$  has super-linear growth if  $\limsup \frac{1}{n} |\Lambda_{n,x}| = \infty$ 

Combined with existing results our result implies:

(Trofimov'84, Lyons-Morris-Schramm'08)

Theorem (Duminil-Copin, G., Raoufi, Severo and Yadin 2018)

Let G be a bounded degree, quasi-transitive graph with superlinear growth, then  $p_c(G) < 1$ .

We say that  $\mathcal{G}$  has spectral dimension at least d > 0 if

$$p_n(x,x) = \mathbb{P}[X_n = x | X_0 = x] \le \frac{c}{n^{d/2}}$$

where X is the simple random walk (SRW) on  $\mathcal{G}$ 

As such spectral dimension is a dynamical property of  ${\cal G}$ 

We say that  $\mathcal{G}$  has spectral dimension at least d > 0 if

$$p_n(x,x) = \mathbb{P}[X_n = x | X_0 = x] \le rac{c}{n^{d/2}}$$
 for all  $x \in \mathcal{G}, n \ge 1$ 

where X is the simple random walk (SRW) on  $\mathcal{G}$ 

For bounded degree graphs isoperimetric dimension > d implies spectral dimension > d (Varopoulos'85)

#### Theorem (Duminil-Copin, G., Raoufi, Severo and Yadin 2018)

Let G a graph of bounded degree with spectral dimension > 4. Then  $p_c(G) < 1$ .

Remark: One noteworthy feature of this result is that it connects percolation threshold, an equilibrium property of  $\mathcal{G}$ , with spectral dimension which is a dynamical property

The Gaussian free field (GFF) on  ${\cal G}$  is a centered Gaussian field

$$\varphi = \{\varphi_x : x \in \mathcal{G}\}$$

Its law  $\mathbb{P}$  is determined by its covariance kernel

$$\mathbb{E}[\varphi_{x}\varphi_{y}] = \frac{g(x,y)}{deg(y)} \sum_{n \ge 0} p_{n}(x,y)$$

g(x, y) is called the Green function of the SRW on  $\mathcal{G}$ 

For any centered Gaussian process  $\psi$  on  $\mathcal G$  with covariances  $\mathcal K(\cdot,\cdot)$ ,

let use define a (bond) percolation process  $\eta_{\mathcal{K}}$  as follows:

#### Definition

Given  $\psi$ , declare an edge xy to be open with probability

$$p_{xy}(\psi) = 1 - \exp(-2(\psi_x + 1)_+(\psi_y + 1)_+)$$

independently of other edges where  $a_{+} = \max(a, 0)$ 

Notice that  $\eta_g$  is a dependent percolation process

#### Proposition

For every  $x \in \mathcal{G}$  one has

$$\mathbb{P}[x \stackrel{\eta_g}{\longleftrightarrow} \infty] \geq \mathbb{E}[\operatorname{sign}(\psi_x + 1)] > 0$$

Let  $\omega_p$  denote the standard percolation on  $\mathcal{G}$  with density p

The main idea is to interpolate between  $\omega_p$  and  $\eta_g$ 

More precisely we prove (here  $g_{\ell}(x, y) = \sum_{\ell \le n \le 2\ell} p_n(x, y)$ )

$$\mathbb{P}[S \xleftarrow{\omega_p \cup \eta_g} T] \leq \mathbb{P}[S \xleftarrow{\omega_{p+1/\ell^2} \cup \eta_g - g_\ell} T]$$

provided  $p_n(x, x)$  decays sufficiently fast

Now iterate this to deduce phase transition

#### Existence of phase transition: open questions

- Prove under the assumption that spectral dimension > 2
- Can we get rid of the assumption of "bounded degree"?

#### Off-critical behavior: truncated two-point function

In the remainder of the talk we will confine ourselves to the hypercubic lattice  $\mathbb{Z}^d$  for  $d \ge 2$ 

The truncated two-point function is of central importance in any model in statistical physics

In the case of percolation it is defined as

#### Off-critical behavior: finite correlation length

It is expected that

$$au_{m{
ho}}(0,x) \sim |x|^{-c} \mathrm{e}^{|x|/\xi(m{
ho})}$$
 as  $x o \infty$ 

where  $\xi(p)$ , called the correlation length, is finite for all  $p \neq p_c$ 

For p close to 0 or 1 (perturbative regime) finiteness of  $\xi(p)$  is not very difficult to show

#### Off-critical behavior: finite correlation length

It is, however, very difficult to prove this for all  $p \neq p_c$ 

In the subcritical regime  $p < p_c$ , this was proved by Menshikov in 1986 and by Aizenman and Barsky in 1987

In the supercritical regime  $p > p_c$  this follows from a famous result by Grimmett and Marstrand in 1990

# Near-critical and critical behavior: correlation length exponent

Infinite correlation length  $\xi(p_c)$  is the hallmark of a critical point

It is expected that  $\xi(p)$  diverges like

$$|p-p_c|^{-
u+o(1)}$$
 as  $p
ightarrow p_c$  for some  $u>0$ 

#### Near-critical and critical behavior: known cases

Currently the correlation length exponent along with other relevant exponents are known rigorously for

- dimension 2 (Schramm, Lawler, Werner, Smirnov)
- ▶ dimension ≥ 11 (Aizenman, Barsky, Hara, Slade, Fitzner, Hofstad)

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### Going beyond independence

Physical and mathematical motivations abound for dependent percolation models

Physically relevant models usually involve interaction

Many models in statistical physics have representations in terms of a percolation process with certain degree of dependence

#### Going beyond independence: zeros of random functions

Geometry of the zero sets of random gaussian functions defined on, e.g.  $\mathbb{R}^n$ , is a classical topic lying at the crossroads between probability theory and geomerty

Two types of gaussian random functions have been studied:

- Random polynomials with (independent) Gaussian coefficients
- Gaussian sum of the eigenfunctions of Laplacian on a compact Riemannian manifold

#### Level-sets of Gaussian Free Field

In short there are many natural percolation models arising from level-sets of gaussian fields with slow decay of correlation

In the discrete set-up a canonical example is the Gaussian free field (GFF) on  $\mathbb{Z}^d$  for  $d \ge 3$ 

Indeed GFF is distributed as a Gaussian sum of the eigenfunctions of (discrete)  $\Delta$  on  $\mathbb{Z}^d$ 

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# GFF: definition

The GFF on  $\mathbb{Z}^d$   $(d \ge 3)$  is a stationary, centered Gaussian field

$$\varphi = \{\varphi_x : x \in \mathbb{Z}^d\}$$

Its law  $\mathbb{P}$  is determined by its covariance kernel

$$\mathbb{E}[\varphi_x \varphi_y] = g(x, y) = E_x[\# \text{visits of SRW to } y]$$

g(x, y) is called the Green function of the SRW on  $\mathbb{Z}^d$ 

#### GFF: correlation function

The green function is asymptotic to the newtonian potential:

$$g(x,y) = g(x-y) \sim |x-y|^{2-d}$$
 as  $|x-y| \to \infty$ 



A (two-dimensional) GFF. (A.Kassel)

#### Level-sets of GFF

We define the level-set above height h as

$$\{\varphi \ge h\} = \{x \in \mathbb{Z}^d : \varphi_x \ge h\}$$

These level-sets form a non-increasing family of site percolation models indexed by height

We can define the corresponding critical value  $h_* = h_*(d)$  as:

$$h_* = \inf\{h \in \mathbb{R} : \mathbb{P}[0 \stackrel{\geq h}{\longleftrightarrow} \infty] = 0\}$$

#### Level-sets of GFF: existence of phase transition



By a soft argument based on the Markov property of GFF it is possible to show that  $h_* \ge 0$  (Bricmont-Lebowitz-Maes'87)

#### Level-sets of GFF: existence of phase transition



It is much more difficult to prove that  $h_* < \infty$ . It was proved for

d = 3 in Bricmont-Lebowitz-Maes'87

#### Level-sets of GFF: existence of phase transition



It was finally proved by Rodriguez and Sznitman (2013) for all

 $d \ge 3$ 

#### Level-sets of GFF: truncated two-point function

Let us recall the definition of truncated two-point function in this context:

$$\tau_h(x,y) = \mathbb{P}\left[ \begin{bmatrix} & \circ & \circ & \circ \\ & \circ & \bullet & \circ \\ & \circ & \bullet & \circ & \circ \\ & \circ & \circ & \circ & \circ \end{bmatrix} = \mathbb{P}[x \xleftarrow{\geq h} y, x \xleftarrow{\geq h} \infty]$$

#### Level-sets of GFF: several possible critical points



It is a very important question to know if  $\bar{h} = h_* = h_{**}$ 

#### Level-sets of GFF: uniqueness of critical point

Theorem (Duminil-Copin, G., Rodriguez, Severo 2019)

For all 
$$d \ge 3$$
,  $\bar{h}(d) = h_*(d) = h_{**}(d)$ 

#### Corollary

For all  $d \ge 3$  and  $h \ne h_*$ , there exists c > 0 and  $\rho \in (0,1]$  such that for all  $x, y \in \mathbb{Z}^d$ ,

$$au_h(x,y) \leq \mathrm{e}^{-c|x-y|^{
ho}}$$

#### Level-sets of GFF: key ideas of the proof

#### Using a sophisticated renormalization argument we first show that



for all  $h \in (\bar{h}, h_{**})$ 

Next we show that these probabilities are superpolynomially close to those for a finitely dependent percolation process

#### Level-sets of GFF: open questions

However using existing techniques we can show that a wide class of finitely dependent percolation processes does not have any such phase except, possibly, at the critical point

This leads to a contradiction since  $(\bar{h}, h_{**})$  can not be nonempty

#### Level-sets of GFF: open questions

- Is ρ actually 1? It seems that the answer might vary based on the dimension (Popov-Texeira'15, Popov-Ráth'15)
- Is it possible to obtain bounds on the exponent of correlation length?
- What happens at the critical point?

# Thank you for your attention!