

Fluctuations in the number of level sets of planar Gaussian fields

Stephen Muirhead (University of Melbourne)

joint work with

Dmitry Belyaev (University of Oxford)

Michael McAuley (University of Helsinki)

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Recall that $N_{LS}(R; \ell)$ and $N_{ES}(R; \ell)$ are the number of level/excursion set components of a Gaussian field f inside a ball $B(R)$, with $c_{LS}(\ell) := \lim_{R \rightarrow \infty} N_{LS}(R; \ell)/(\pi R^2)$ and c_{ES} similar.

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Theorem (Belyaev, McAuley, M., '19)

For the BF field

$$\text{Var}[N_{LS}(R; \ell)] \gtrsim R^2$$

for all $\ell \in \mathbb{R}$ such that $c'_{LS}(\ell) \neq 0$, and similarly for $N_{ES}(R; \ell)$.

For the RPW

$$\text{Var}[N_{LS}(R; \ell)] \gtrsim R^3$$

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The orders R^2/R^3 for the BF/RPW are quite natural and we expect them to be tight (at least for generic levels).

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Corollary (Belyaev, McAuley, M.)

The variance of N_{LS}/N_{ES} is of 'full order' for the following levels:

- ▶ *BF field ('full order' = R^2)*
 - ▶ *LS: $\ell \in (-\infty, -1.38) \cup (1.38, \infty)$*
 - ▶ *ES: $\ell \in (-\varepsilon, 0.64) \cup (1.02, \infty)$*
- ▶ *RPW ('full order' = R^3)*
 - ▶ *LS: $\ell \in (-\infty, -1] \cup [1, \infty)$*
 - ▶ *ES: $\ell \in (-\infty, 0) \cup (0, 0.87) \cup [1, \infty)$*

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- 2) Gaussian fields can 'breathe', i.e. f and $f + \varepsilon$ are statistically indistinguishable on a compact domain $D \subset \mathbb{R}^2$ for small enough $\varepsilon > 0$ (depending on D).

There are three main ideas in the proof:

- 1) One can bound global topological events by local observables.
- 2) Gaussian fields can 'breathe', i.e. f and $f + \varepsilon$ are statistically indistinguishable on a compact domain $D \subset \mathbb{R}^2$ for small enough $\varepsilon > 0$ (depending on D).
- 3) By a coupling argument, one can deduce variance lower bounds for the number of level sets at a fixed level by considering the **change as the level varies**.

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Hence, by the Kac-Rice formula

$$\begin{aligned} \mathbb{E}[N_{LS}(R; \ell)^2] &\leq \mathbb{E}[N_{LS}(R; \ell)(N_{LS}(R; \ell) - 1)] \\ &= \int_{x, y \in B(R)} \mathbb{E}[|\det(\nabla^2 f(x) \nabla^2 f(y))| \mid \nabla f(x) = \nabla f(y) = 0] \\ &\quad \times \varphi_{\nabla f(x), \nabla f(y)}(0, 0) \, dx dy \end{aligned}$$

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We need a slightly refined version valid for small height windows:

Lemma (M. 20)

For every $\ell \in \mathbb{R}$, $R > 0$ and $a < b$,

$$|N_{LS}(R; b) - N_{LS}(R; a)| \lesssim \# \text{ critical points in } B(R) \text{ with level in } [a, b].$$

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2) Gaussian fields can 'breathe'

We next consider how small $\varepsilon > 0$ needs to be, as a function of R , such that the fields

$$f|_{B(R)} \quad \text{and} \quad (f + \varepsilon)|_{B(R)}$$

are close in total variation distance (and hence any functionals of the fields are also close in TV-distance).

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Nevertheless, an **approximate** version of this comparison is true and this is sufficient for our purposes.

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Fact 1: Since the covariance kernel K is in $L^1(\mathbb{R}^2)$, f has a 'moving average representation'

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Fact 2: Let $W_D = W\mathbb{1}_D$ be white noise on a compact domain $D \subset \mathbb{R}^2$. Then $f_D = q \star W_D$ has an 'orthogonal decomposition'

$$f_D = q \star W_D = \sum_{i \geq 1} Z_i(q \star \varphi_i)$$

where Z_i is a sequence of independent standard Gaussians, and φ_i is any orthonormal basis of $L^2(D)$.

Fact 3: For a standard Gaussian vector $Z = (Z_1, \dots, Z_n)$,

$$d_{TV}(Z, Z + \varepsilon) \lesssim \varepsilon\sqrt{n}.$$

Proof. Use Pinsker's inequality $d_{TV}(\mu, \nu) \lesssim \sqrt{d_{KL}(\mu||\nu)}$, and then the additivity of relative entropy for product measures.

Putting these together we have:

Proposition

For the BF field

$$d_{TV}(f|_{B(R)}, (f + \varepsilon g)|_{B(R)}) \lesssim \varepsilon R$$

where $g = (q \star \mathbb{1}_{B(2R)})|_{B(R)} \approx (\int q) \times \mathbb{1}_{B(R)}$.

Proof. Divide the plane into unit boxes D_i , then decompose orthogonally each $W\mathbb{1}_{D_i}$ with $\varphi_1 = c\mathbb{1}_{D_i}$, and then shift the Gaussians Z_1 by ε in each box.

Upshot: Since $\int q > 0$ for the BF field, the number of level/excursion set components of

$$f|_{B(R)} \quad \text{and} \quad (f + \varepsilon)|_{B(R)}$$

are close in total variation distance as soon as $\varepsilon \ll 1/R$.

The proceeding argument works for any field f whose spectral density ρ does not vanish at the origin (the condition $\int q > 0$ is equivalent to $\rho(0) > 0$).

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However, for the RPW this argument does not work ($K \notin L^1$, and the spectral measure is Lebesgue on the unit circle).

There is a good reason for this – the RKHS of the RPW consists of solutions of the Helmholtz equation $\nabla f = -f$ which does not contain constant functions (or any approximation of them).

So one cannot expect f and $f + \varepsilon$ to be comparable in total variation distance.

For the RPW we use a slightly different approach. Recall the orthogonal decomposition of the RPW

$$f(x) = f(r, \theta) = \Re\left(\sum_{n=-\infty}^{n=\infty} Z_n e^{2\pi i n \theta} J_{|n|}(r)\right)$$

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Proof. Again use Pinsker's inequality.

Upshot: Fix $\ell \neq 0$. Then the fields

$$f_n|_{B(R)} \quad \text{and} \quad \left(f_n \times \frac{\ell + \varepsilon}{\ell}\right)|_{B(R)}$$

are close in total variation distance as soon as $\varepsilon \ll 1/\sqrt{R}$.

Hence so are the number of components of the level sets

$$\{f|_{B(R)} = \ell\} \quad \text{and} \quad \{f|_{B(R)} = \ell + \varepsilon\}.$$

To sum-up, we have:

BF: The number of level/excursion set components of $f|_{B(R)}$ at levels

$$l \quad \text{and} \quad l + \varepsilon$$

are close in total variation distance for $\varepsilon \ll 1/R$.

RPW: The number of level/excursion set components of $f|_{B(R)}$ at levels

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are close in total variation distance for $\varepsilon \ll 1/\sqrt{R}$.

The fact that we can shift levels by $1/\sqrt{R}$ ($\gg 1/R$) for the RPW is a manifestation of the **strong degeneracies** in the RPW.

3) A coupling method for variance lower bounds

Recently Chatterjee proposed a general coupling method for proving lower bounds on variances:

Lemma (Chatterjee 2017)

Let X and Y be two random variables defined on the same probability space. Then for every $a \leq b$,

$$P(a \leq X \leq b) \leq \frac{1}{2} \left(1 + P(|X - Y| \leq b - a) + d_{TV}(X, Y) \right).$$

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Hence, if X_n is a sequence of random variables and there exist another sequence Y_n , and constants $\sigma_n, \delta > 0$, such that:

- ▶ $P(|X_n - Y_n| > \delta \sigma_n) > 1/4$
- ▶ $d_{TV}(X_n, Y_n) < 1/8$.

then X_n fluctuates on the scale $\gtrsim \sigma_n$, and hence $\text{Var}(X_n) \gtrsim \sigma_n^2$.

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$$c_{ES}(\ell) := \lim_{R \rightarrow \infty} X_R / (\pi R^2).$$

By controlling the error in the above convergence, if $c'_{ES}(\ell) \neq 0$ we can prove that

$$\mathbb{E}[Y_R - X_R] = \varepsilon_R R^2 c'_{ES}(\ell) + o(\varepsilon_R R^2).$$

Moreover, by considering critical points with heights in $[\ell, \ell + \varepsilon_R]$, we have

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To conclude, for the BF we have $d_{TV}(X_R, Y_R) \lesssim 1$ as long as $\varepsilon_R \lesssim 1/R$, and so the coupling method gives $\text{Var}[X_R] \gtrsim (\varepsilon_R R^2)^2 \gtrsim R^2$.

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On the other hand, if $\ell \neq 0$ then for the RPW we have $d_{TV}(X_R, Y_R) \lesssim 1$ as long as $\varepsilon_R \lesssim 1/\sqrt{R}$, and so the coupling method gives $\text{Var}[X_R] \gtrsim (\varepsilon_R R^2)^2 \gtrsim R^3$.

Validating the conditions

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Recall that we can prove

$$c_{ES}(\ell) = \int_{\ell}^{\infty} p_{m^+}(x) - p_{s^-}(x) dx = \int_{\ell}^{\infty} p_{m^+}(x) - p_s(x)g(x) dx.$$

where $p_s(\ell)$ is the density of saddle points, and

$g(\ell) := \mathbb{P}[0 \text{ is a lower connected saddle} \mid 0 \text{ is a saddle with height } \ell]$.

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Hence $c'_{ES}(\ell) \neq 0$ is equivalent to

$$g(\ell) \neq p_{m^+}(\ell)/p_s(\ell)$$

and the right-hand side is explicitly computable.

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This is surprisingly hard to prove in general. But by considering the field conditioned to have a saddle at 0 at height x (and using explicit properties of the BF/RPW) we can prove:

Lemma

For both the BF and RPW, $g(\ell)$ is strictly increasing in ℓ .

Since $g(0) = 1/2$ (by symmetry), this yields

$$g(\ell) = \begin{cases} > 1/2 & \text{if } \ell > 0, \\ < 1/2 & \text{if } \ell < 0, \end{cases}$$

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which gives regions where $g(\ell) \neq p_{m^+}(\ell)/p_s(\ell)$ and so $c'_{ES}(\ell) \neq 0$.

A similar argument works for $c'_{LS}(\ell)$.

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2) For the RPW, there are no local maxima at levels $\ell < 0$.

Hence $c'_{ES}(\ell) > 0$ for all $\ell \leq 0$.

The result of this argument is explicit, but somewhat modest:

Corollary (Belyaev, McAuley, M.)

We have $c'_{LS}(\ell) \neq 0$ and $c'_{ES}(0) \neq 0$ (and hence 'full order' variance, except at $\ell = 0$ for the RPW) for the following levels:

▶ *BF field*

▶ *LS: $\ell \in (-\infty, -1.38) \cup (1.38, \infty)$*

▶ *ES: $\ell \in (-\varepsilon, 0.64) \cup (1.02, \infty)$*

▶ *RPW*

▶ *LS: $\ell \in (-\infty, -1] \cup [1, \infty)$*

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1. We believe that c_{LS} (resp. c_{ES}) has at most two (resp. one) critical points. Is this true? Can we at least show that the set of critical points is finite?
2. Are fluctuations of the number of level/excursion sets genuinely of lower order if $c'_{LS}(\ell) = 0$ and $c'_{ES}(\ell) = 0$? We think this is **not** true for the BF field, but might be true for the RPW (by analogy with the length of the nodal set, where this phenomena occurs and is known as **Berry cancellation**). If it's true, how small are the fluctuations in these cases?

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3. Matching upper bounds / leading order constants / CLTs?

Thank you!

References:

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