# A beginner's introduction to the Liouville quantum gravity metric

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# Liouville quantum gravity (LQG)

- Model for a random two-dimensional Riemannian manifold.
- A  $\gamma$ -LQG surface parametrized by  $\mathbb{C}$  informally is a random Riemannian manifold with area measure  $e^{\gamma h} d^2 z$  and the Riemannian metric tensor  $e^{2\xi h} (dx^2 + dy^2)$  for some parameter  $\gamma$ and some  $\xi = \xi(\gamma)$  where h is the Gaussian free field on  $\mathbb{C}$ .
- Considering the scaling exponent relating the the length and area suggests  $\xi$  and  $\gamma$  should be related by  $\xi = \gamma/d_{\gamma}$  is the "dimension" of the surface.
- Because of the roughness of *h*, LQG surfaces cannot be Riemannian manifolds in the usual sense, defining these rigorously has been a major challenge.

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# Ranges of $\gamma$ and $\xi$

- Usual parameter range:  $\gamma \in (0, 2)$  (subcritical).
- $\xi$  is an increasing function of  $\gamma$  on (0, 2] and  $\xi(2) \approx 0.41$ .
- $\xi = \xi(2)$  (critical) and  $\xi > \xi(2)$  (supercritical) corresponds to  $\gamma = 2$  and  $\gamma \in \mathbb{C}$  with  $|\gamma| = 2$  respectively.
- We shall only discuss the subcritical case.

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# Motivation and history-some physics buzzwords

- LQG surfaces were first considered by Polyakov in 1981-motivation came from string theory.
- Also connected to quantum field theory.
- LQG surface parametrized by a domain is "uniform sample from the space of Riemannian metric tensors weighted by  $(\det \Delta_g)^{-c_M/2}$ " where  $\Delta_g$  is the Laplace-Beltrami operator and  $c_M$ is the matter central charge.

• 
$$c_M = 25 - 6Q^2$$
 where

$$Q = \left(\frac{2}{\gamma} + \frac{\gamma}{2}\right)$$

•  $c_M \in (-\infty, 1]$  corresponds to subcritical and critical cases, whereas  $c_M \in (1, 25)$  corresponds to the supercritical regime.

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## Motivation-random planar maps

- Uniform planar maps, i.e., uniform planar triangulation are believed converge to LQG with  $c_M = 0$ , i.e.,  $\gamma = \sqrt{8/3}$ .
- For other values of γ ∈ (0, 2), γ-LQG is believed to be the scaling limit of planar maps decorated with statistical physics models.
- For example, a planar map decorated with spanning tree is believed to converge to  $\gamma$ -LQG for  $\gamma = \sqrt{2}$ .

Gaussian free fields (GFF)

• For a bounded open set U, and a compactly supported smooth function  $f \in C_c^{\infty}(U)$  consider the Dirichlet norm of f given by

$$||f||_{\nabla}^2 = \int_U |\nabla f|^2 d^2 z.$$

- Let  $H_0^1(U)$  denote the Hilbert space closure of  $C_c^{\infty}(U)$  with respect to the Dirichlet norm.
- For an orthonormal basis  $f_1, f_2, \ldots$ , of  $H_0^1(U)$ , zero boundary GFF is formally defined as

$$h = \sum_{i} X_i f_i$$

where  $X_i$  are i.i.d. N(0, 1).

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Gaussian free fields (GFF)

- Does not make sense as a random function.
- We can make sense of this as a random generalised function (distribution) on U.
- We set

$$(h,\phi)_{\nabla} = \sum_{i=1}^{\infty} X_i(f_i,\phi)_{\nabla}$$

which is distributed as  $N(0, \|\phi\|_{\nabla}^2)$ .

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## Zero boundary GFF

• One can then define the action of h on  $\phi \in C_c^{\infty}(U)$  by setting

$$(h,\phi) = -2\pi(h,\Delta^{-1}\phi)_{\nabla}$$

where  $\Delta^{-1}$  is the inverse Laplacian with zero boundary conditions.

- *h* is thus defined as a random element of the continuous dual of  $C_c^{\infty}(U)$ .
- Zero boundary GFF is conformally invariant: if  $\Phi : U \to V$  is a conformal map, and h is a zero boundary GFF on V, then  $h \circ \Phi$  is a zero boundary GFF on U.

# Whole plane GFF

- On C, constant functions have zero Dirichlet energy, it does not induce a norm on H<sup>1</sup><sub>0</sub>(C).
- One defines the whole plane GFF h similarly as before as a random element of  $(C_c^{\infty}(\mathbb{C}))'$  modulo constants.
- To make sense of h as an element of the dual itself we shall fix the normalization

$$\mathbf{Av}(h,\mathbb{T})=0$$

where  $\mathbf{Av}(h, \mathbb{T})$  denotes the GFF integrated against the uniform measure on the unit circle (one can show that this can be defined).

• With this normalization h is no longer conformally invariant. However, using conformal invariance at the level of a generalized function modulo constants, we get

$$h(r\cdot) - \mathbf{Av}(h, \mathbb{T}_r) \stackrel{d}{=} h(\cdot).$$

Key properties of GFF

Domain Markov property

Roughly speaking, on a domain U, one can write

$$h = h_1 + h_2$$

where  $h_1$  is the harmonic extension of the boundary condition on  $\partial U$ and  $h_2$  is an independent GFF on U with zero boundary condition.

Circle average process

$$B(t) = \mathbf{Av}(h, \mathbb{T}_{e^t})$$

is distributed as a standard two-sided Brownian motion.

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## Mollification of GFF

- To define LQG measure and metric rigorously, one considers a mollification of GFF which is a random function and then takes an appropriate limit for the mollification parameter.
- One particular convenient choice of mollification is to take convolution with the heat kernel.

$$h_{\varepsilon}(z) = \int_{\mathbb{C}} h(w) \left(\frac{1}{\pi \varepsilon^2} e^{-|z-w|^2/\varepsilon^2}\right) d^2 w.$$

• It is possible to consider other mollification schemes.

#### Renormalisation and LQG area measure

• Consider the family of random measures on  $\mathbb{C}$ :

$$\mu_{h,\varepsilon} = \varepsilon^{\gamma^2/2} e^{\gamma h_{\varepsilon}(z)} d^2 z.$$

- For  $\gamma \in (0, 2)$ , one can show that  $\mu_{h,\varepsilon}$  converges to a random measure  $\mu_h$  on  $\mathbb{C}$  almost surely in the randomness of h.
- This limit is the LQG area measure.
- Duplantier-Sheffield (2011), Rhodes-Vargas (2014).

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#### Renormalisation and LQG metric

• For  $\xi > 0$ , define the following metric induced by  $h_{\varepsilon}$ :

$$D_{h,\varepsilon}(z,w) = \inf_{\pi} \int_0^1 e^{\xi h_{\varepsilon}(\pi(t))} |\pi'(t)| dt$$

where the infimum is taken over all sufficiently nice and properly parametrized paths between z and w.

• It is shown that there exists a family of constants  $a_{\varepsilon}$  such that

$$a_{\varepsilon} = \varepsilon^{1 - \xi Q + o(1)}$$

as  $\varepsilon \to 0$  such that for  $\xi < \xi(2) = \frac{2}{d_2}$  the metrics

$$\tilde{D}_{h,\varepsilon} = a_{\varepsilon}^{-1} D_{h,\varepsilon}$$

converges in probability to a metric  $D_h$ .

# Renormalisation and LQG metric

- Ding-Gwynne (2018): showed the existence of  $d_{\gamma}$  via discrete approximation schemes.
- Ding-Dubedat-Dunlop-Falconet (2020): showed the tightness of the renormalised metrics.
- Gwynne-Miller (2021): established uniqueness.

## Basic properties of LQG metric

- $\mathbb{C}$ , equipped with  $D_h$  is a length space.
- $d_{\gamma}$  is the Hausdorff dimension of  $\mathbb{C}$  equipped with the metric  $D_h$ .
- $D_h$  is local, i.e., for any open set U, the internal metric  $D_h(\cdot, \cdot; U)$  is determined by the restriction of h to U.
- Induces the Euclidean topology on C and is bi-Hölder with respect to the Euclidean metric (quantitative estimates available).

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#### Basic properties of LQG metric

- One can similarly define the LQG metric  $D_{h}$  where h = h + f, GFF plus a continuous function.
- Weyl Scaling:  $D_{h+f} = e^{\xi f} \cdot D_h$  almost surely.
- Translation invariance: for each  $z \in \mathbb{C}$ ,  $D_{h(\cdot+z)} = D_h(\cdot+z, \cdot+z)$  almost surely.
- Coordinate change formula: almost surely for any fixed r > 0 and all  $z, w \in \mathbb{C}$ ,

$$D_{h(r\cdot)}(rz, rw) = r^{\xi Q} D_{h(\cdot)}(z, w).$$

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# Geodesics in LQG

- Almost surely, for any two  $z, w \in \mathbb{C}$  there exists a  $D_h$  geodesic between z and w. For fixed z and w the geodesic is almost surely unique, we shall denote it by  $\Gamma(z, w)$ . Geodesics will always be parametrized by the LQG distance.
- For each  $z \in \mathbb{C}$ , almost surely there exists a unique infinite geodesic started from z.
- Scale invariance of the geodesic: For the infinite geodesic  $\Gamma$  started from the origin we have

$$(h, \Gamma_{\cdot}) = (h(r \cdot) - \mathbf{Av}(h, \mathbb{T}_r), r^{-1} \Gamma_{r \in Q_e \in \mathbf{Av}(h, \mathbb{T}_r)})$$

in distribution.

- Almost surely there are no bigeodesics.
- Gwynne-Pfeffer-Sheffield (2020).

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#### Coalescence of geodesics

• Fix  $z \in \mathbb{C}$ . For any s > 0, there is t > 0 such that all geodesics from z to points at LQG distance larger than s coincide within the LQG metric ball centred at z with radius t.



#### Coalescence of geodesics

• Let  $\mathbf{Coal}_{r,K}$  denote the event

$$\mathbf{Coal}_{r,K} = \bigg\{ \cap_{z \in \mathbb{T}_r, w \in \mathbb{T}_{Kr}} \Gamma(z, w) \neq \emptyset \bigg\}.$$

• For K sufficiently large,  $\mathbb{P}(\mathbf{Coal}_{r,K})$  is bounded away from 0 uniformly in r.



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# Critical and supercritical LQG metrics

- Many similar results have very recently been proved for critical and supercritical LQG metrics as well.
- The critical LQG metric still induces the Euclidean topology but is no longer bi-Holder with respect to the Euclidean metric.
- Supercritical LQG metric has Hausdorff dimension  $\infty$ .

#### References

- Nathanaël Berestycki and Ellen Powell. Gaussian free field, Liouville quantum gravity and Gaussian multiplicative chaos.
- Ewain Gwynne, Nina Holden, and Xin Sun. Mating of trees for random planar maps and Liouville quantum gravity.
- Jian Ding, Julian Dubedat and Ewain Gwynne. An introduction to the Liouville quantum gravity metric.

## Thank you

Questions?

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