

# Some non-existence results for the stochastic wave equation

The effect of noise on blow-up of deterministic systems

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## Some References

- Julian Fernández Bonder and Pablo Groisman. Time-space white noise eliminates global solutions in reaction-diffusion equations. *Phys. D*, 238(2):209-215, 2009.
- Mohammud Foondun and Eulalia Nualart. The Osgood condition for stochastic partial differential equations. to appear in *Bernouilli*, 2020.
- Robert Glassey. Blow-up theorems for nonlinear wave equations. *Math. Z.*, 132:183-203, 1973.
- Mohammud Foondun and Eulalia Nualart. Non-existence results for stochastic wave equations in one dimension

# The heat equation problem

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- Consider the following ODE

$$\frac{dy}{dt} = b(y), \quad y(0) = y.$$

- The solution blows up in finite time if and only if

$$T = \int_y^{\infty} \frac{1}{b(s)} ds$$

is finite.

- We now consider the following SDE:

$$X_t = X_0 + \int_0^t b(X_s) ds + B_t.$$

- The solution blows up almost surely if and only if

$$\int_0^\infty \frac{1}{b(s)} ds < \infty.$$

- This is a special case of Feller's test for explosion.
- The intuition is that the Brownian motion pushes the solution up so that the deterministic part makes the solution blow up.

# The deterministic heat equation

- We now consider the heat equation with Dirichlet boundary condition.

$$\frac{\partial u}{\partial t} = \Delta u + b(u) \quad \text{on } [0, 1]$$

- The solution will blow-up in finite time if

$$\int_{.}^{\infty} \frac{1}{b(s)} ds < \infty$$

and the initial condition  $u_0$  is large enough.

- The point is that the presence of the Dirichlet Laplacian makes it that the above integral condition is not enough for blow-up.

# The stochastic heat equation

- We now look at the following stochastic heat equation

$$\frac{\partial u}{\partial t} = \Delta u + b(u) + \dot{W} \quad \text{on } [0, 1]$$

- It turns out since the noise terms pushes the solution up, the following condition guarantees blow-up no matter what the initial condition is

$$\int_0^\infty \frac{1}{b(s)} ds < \infty.$$

- Question: What kind of condition do we have for the stochastic wave equation?



# Proof ideas

- Look at  $y(t) := \int_0^1 u(t, x)\phi(x)dx$
- Use comparison arguments.
- For the deterministic case, compare with a non-linear ODE which can be analysed
- For the stochastic case, compare with an SDE which blows up according to Feller's test.
- If  $y(t)$  blows up, then  $\sup_{x \in [0,1]} |u(t, x)|$  blows up as well.

# The wave equation

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- Consider the second order ODE:

$$\frac{d^2y}{dt^2} = b(y) \quad y(0) = \alpha, \quad y'(0) = \beta.$$

- This is equivalent to

$$y(t) = \alpha + \beta t + \int_0^t (t-s)b(y(s)) ds$$

- The solution blows up if and only if

$$\int_{\alpha}^{\infty} \frac{1}{[\beta^2 + 2 \int_{\alpha}^s b(r) dr]^{1/2}} ds < \infty$$

- What is the role of this condition on blow-up properties of stochastic equations?

# The proof

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$$y''(t)y'(t) = b(y(t))y'(t) \quad t \geq 0$$

- This is equivalent to

$$y'(t)^2 - y'(0)^2 = 2 \int_0^t b(y(s)) dy(s) = 2 \int_{\alpha}^{y(t)} b(r) dr.$$

- This is of the form  $y'(t) = F(y(t))$ . So that the integral condition for the first order ODE applies and we obtain the required condition.

# The deterministic wave equation

- Consider the following wave equation with Dirichlet boundary condition.

$$\frac{\partial^2 u}{\partial t^2} = \Delta u + b(u) \quad \text{on } [0, 1]$$

- Set

$$\alpha = \int_0^1 \phi(x)u(x, 0)dx \quad \beta = \int_0^1 \phi(x)u_t(x, 0)dx$$

- The solution blows up in finite time provided that

$$T = \int_{\alpha}^{\infty} \left[ \mu\alpha^2 + \beta^2 - \mu s^2 + 2 \int_{\alpha}^s b(r)dr \right]^{-1/2} ds < \infty$$

# The Stochastic wave equation

- We look at

$$\frac{\partial^2 u}{\partial t^2} = \Delta u + b(u) + \dot{W} \quad \text{on } [0, 1]$$

with Dirichlet boundary conditions.

- Suppose that for  $\alpha, \beta > 0$ , we have

$$T(\alpha, \beta) := \int_{\alpha}^{\infty} \frac{1}{[\beta^2 + 2 \int_{\alpha}^s b(r) dr]^{1/2}} ds < \infty$$

- Then the solution blows up with a positive probability.

## The idea behind the proof

- We look at the integral formulation of the solution.

$$\begin{aligned}u(t,x) &= \int_0^1 G(t,x,y)v_0(y) dy + \frac{\partial}{\partial t} \left( \int_0^1 G(t,x,y)u_0(y) dy \right) \\ &+ \int_0^t \int_0^1 G(t-s,x,y) W(ds dy) \\ &+ \int_0^t \int_0^1 G(t-s,x,y)b(u(s,y)) ds dy \quad \text{a.s.}\end{aligned}$$

- A major difficulty is that we do not have a comparison principle for the wave equation.
- But an important observation is that for small times we can still use some kind of comparison argument together with the fact that the stochastic part gets large with a positive probability.

- The kernel  $G(t, x, y)$  is given by the following

$$G(t, x, y) := \sum_{n=1}^{\infty} \frac{\sin(n\pi t)}{n\pi} \varphi_n(x) \varphi_n(y),$$

where  $\varphi_n(x) = \sqrt{2} \sin(n\pi x)$ ,  $n \geq 1$ .

- We look at  $y(t) = \int_0^1 u(t, x) \phi_1(x) dx$ .
- The idea is to compare  $y(t)$  with an integral equation for a certain time interval only.
- The integral condition will then follow.



- The stochastic part is given by

$$\int_0^t \int_0^1 \int_0^1 G(t-s, x, y) \varphi_1(x) W(ds dy) dx,$$

- This can be rewritten as

$$\begin{aligned} M(t) &= \int_0^t \int_0^1 \sin(\pi(t-s)) \varphi_1(y) W(dy ds) \\ &= C \int_0^t \sin(\pi(t-s)) dB_s \\ &= C\pi \int_0^t \cos(\pi(t-s)) B_s ds. \end{aligned}$$

- Using this and the support theorem for Brownian motion, we can show that in a certain time interval the above quantity can get large with a positive probability.

- The function  $y(t)$  can then be compared to an integral equation of the form

$$y(t) = A + B(t - t_0) + 2 \int_{t_0}^t (t - s)b(y(s)) ds + L \quad t \in [t_0, T].$$

- For blow-up we need

$$\int_{A+L}^{\infty} \frac{1}{[B^2 + 2 \int_{A+L}^s b(r) dr]^{1/2}} ds < \infty,$$

- Consider

$$y(t) = A + Bt + \int_0^t (t-s)b(y(s)) ds + G(t) \quad t \in [0, T].$$

- $G(t)$  grows to infinity as  $t \rightarrow \infty$
- For blow-up we need

$$\int_{\alpha}^{\infty} \frac{1}{[\beta^2 + 2 \int_{\alpha}^s b(r) dr]^{1/2}} ds < \infty,$$

# The wave equation on the whole line

- We look at

$$\frac{\partial^2 u}{\partial t^2} = \Delta u + b(u) + \dot{W}$$

on the whole line

- Suppose that for  $\alpha, \beta > 0$ , we have

$$T(\alpha, \beta) := \int_{\alpha}^{\infty} \frac{1}{[\beta^2 + 2 \int_{\alpha}^s b(r) dr]^{1/2}} ds < \infty$$

- Then the solution blows up almost surely.

## The idea behind the proof

- We look at the integral formulation of the solution.

$$\begin{aligned}u(t,x) &= \int_{-\infty}^{\infty} G(t,x,y)v_0(y) dy + \frac{\partial}{\partial t} \left( \int_{-\infty}^{\infty} G(t,x,y)u_0(y) dy \right) \\ &+ \int_{-\infty}^{\infty} \int_0^1 G(t-s,x,y) W(ds dy) \\ &+ \int_{-\infty}^{\infty} \int_0^t G(t-s,x,y)b(u(s,y)) ds dy \quad \text{a.s.}\end{aligned}$$

- The following is a Gaussian process

$$g(t,x) := \int_0^t \int_{-\infty}^{\infty} G(t-s,x-y)W(dy ds).$$

- For fixed  $x \in \mathbf{R}$ , almost surely,

$$\limsup_{t \rightarrow \infty} \frac{g(t, x)}{t \sqrt{\log \log t}} > 1.$$

- This requires a bit of Gaussian theory to prove.
- We can then compare the mild solution to an integral equation which blows up.